

HYBRID SYSTEMS FOR QUANTUM INFORMATION

C H R I S W I L S O N

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NICOLAS SANGOUARD, MIKAEL AFZELIUS

CHALMERS

MC2
Microtechnology and Nanoscience

HYBRID QUANTUM NETWORK

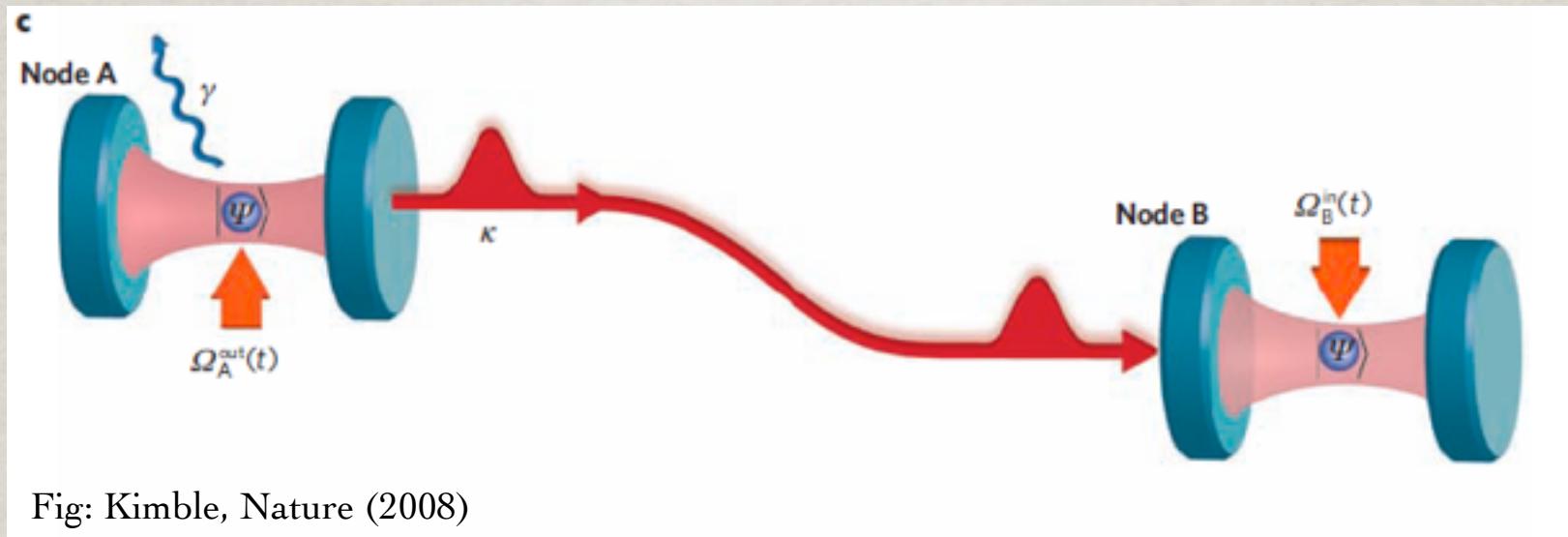


Fig: Kimble, Nature (2008)

- ✿ Telecom photons to distribute quantum information
- ✿ Superconducting circuits to process at nodes

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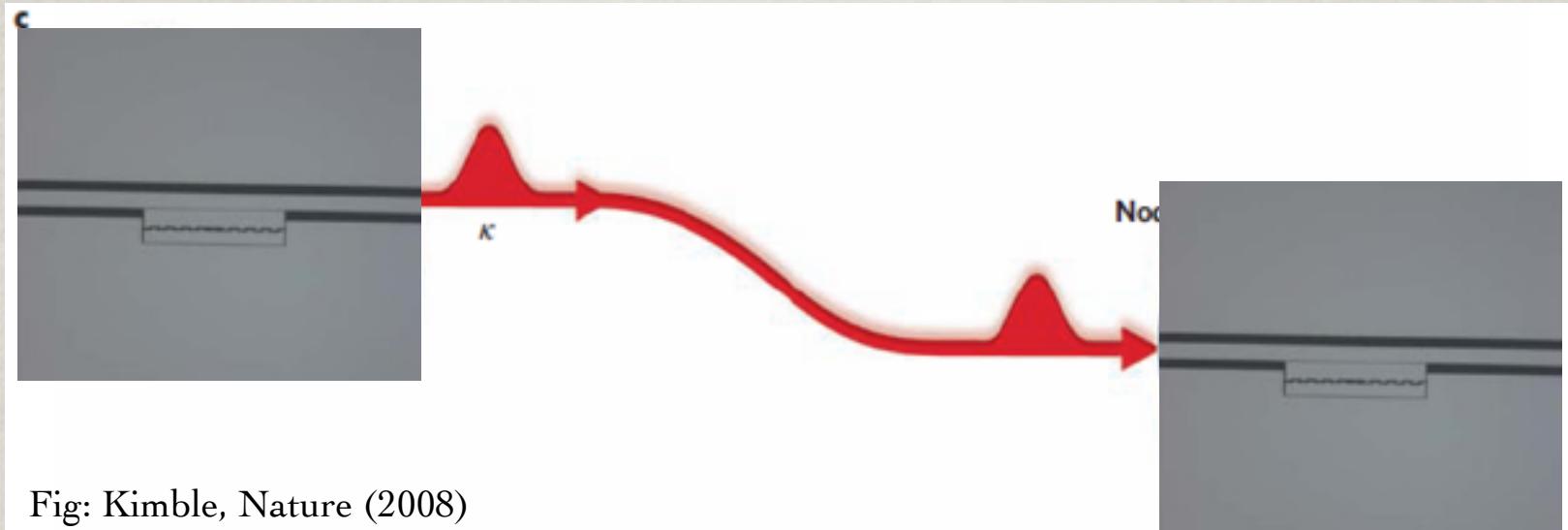


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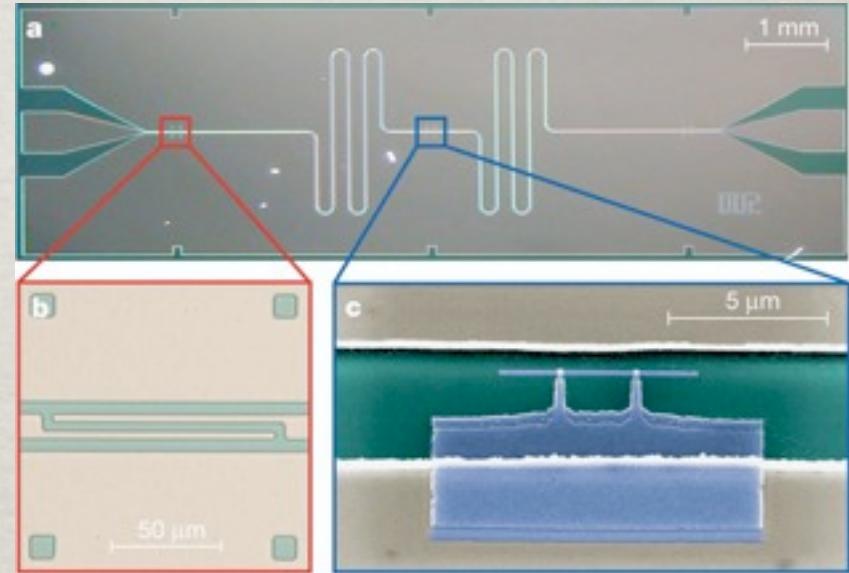
- ✿ Telecom photons to distribute quantum information
- ✿ Superconducting circuits to process at nodes

QUANTUM CIRCUITS AND ARTIFICIAL ATOMS

CIRCUIT QUANTUM ELECTRODYNAMICS

Atoms \Rightarrow Qubits
3D Cavity \Rightarrow 1D on-chip resonator

- In 2004, Wallraff et. al. introduced circuit QED
- Coupling strengths not accessible in cavity QED are possible.
- “Atoms” can be engineered.



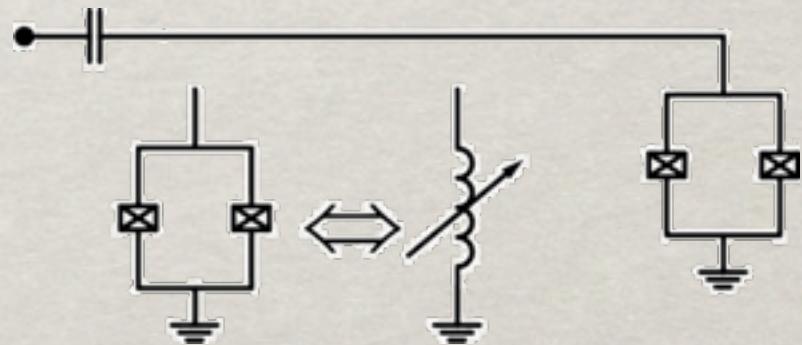
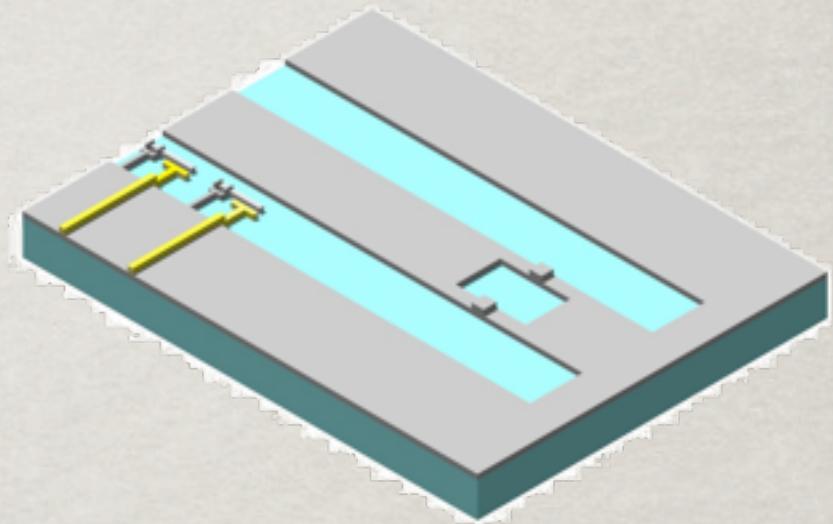
Wallraff et. al. Nature, 2004.

CIRCUIT QUANTUM ELECTRODYNAMICS

- ✿ Vacuum Rabi splitting. Wallraff, Nature, 2004.
- ✿ Photon number distribution. Schuster, Nature, 2007.
- ✿ Entanglement between two qubits. Majer, Nature, 2007.
- ✿ Resolving Vacuum fluctuations. Fragner, Science, 2008.
- ✿ Generation and resolving of number states. Hoffheinz, Nature, 2008,2009.
- ✿ Simple quantum algorithms. DiCarlo Nature, 2009.
- ✿ And more...

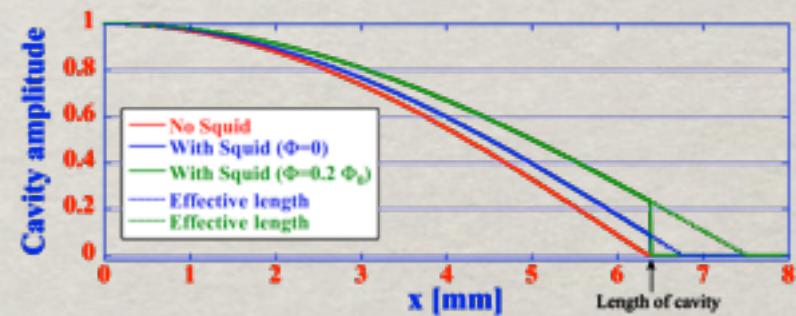
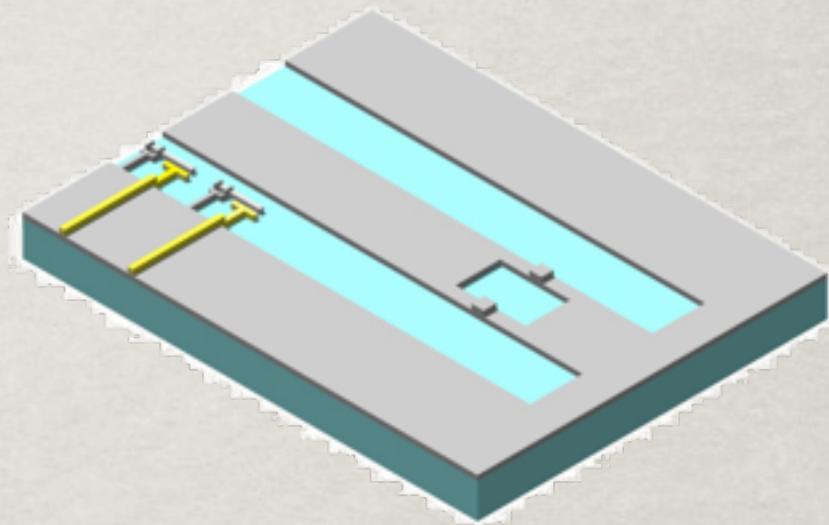
TUNABLE TRANSMISSION LINE RESONATOR

- Transmission line terminated in a tunable impedance.
- Tunable impedance: the Josephson inductance of a SQUID.
- M. Wallquist et al. constructed a control-phase gate protocol



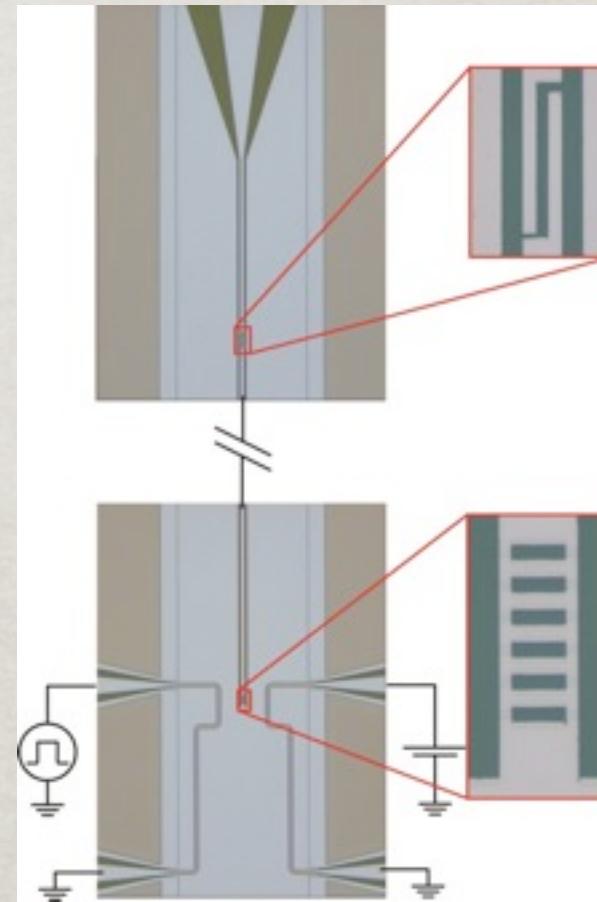
TUNABLE TRANSMISSION LINE RESONATOR

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DEVICE FABRICATION

Transmission line and SQUID
fabricated in Al. Tunnel junctions
fabricated from native Al_2O_3 .

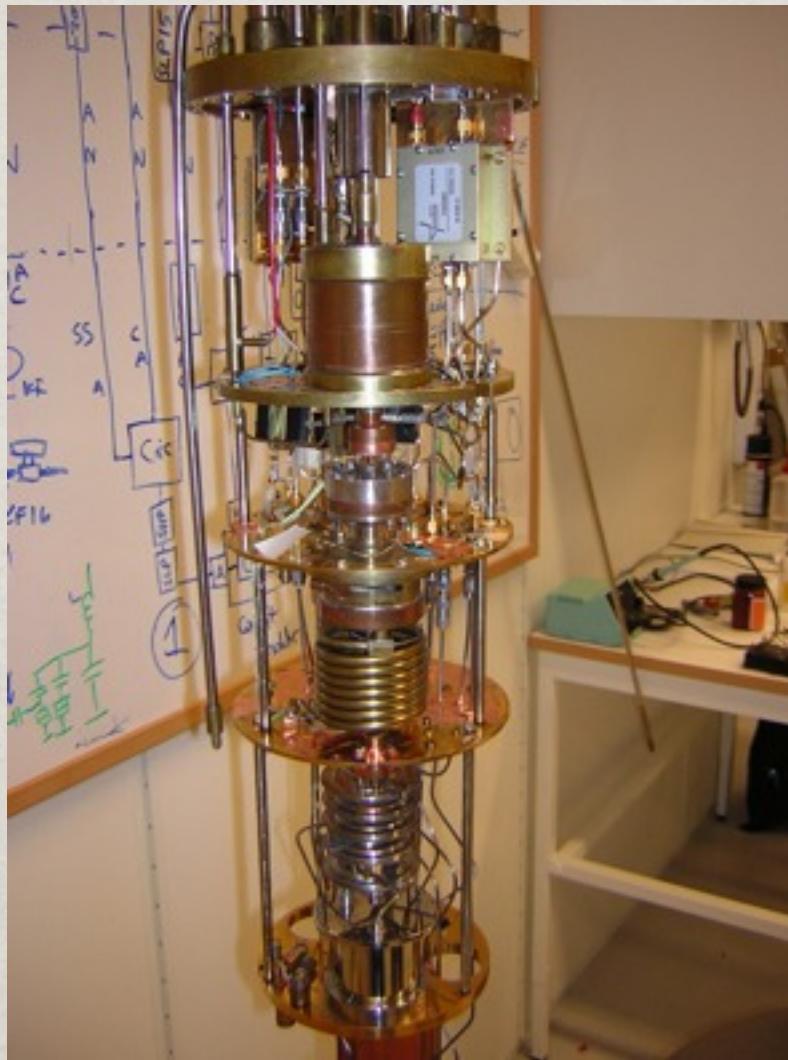


Coupling
capacitance

Array of
SQUIDs

On-chip tuning lines for fast tuning

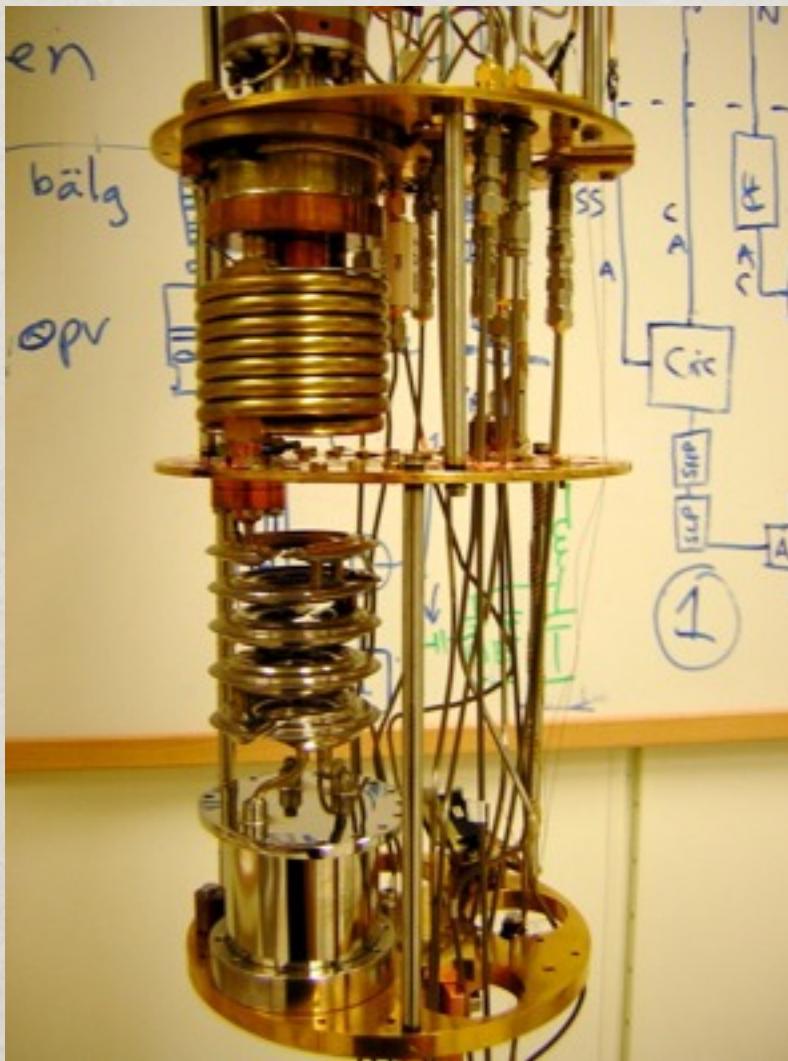
DILUTION FRIDGE



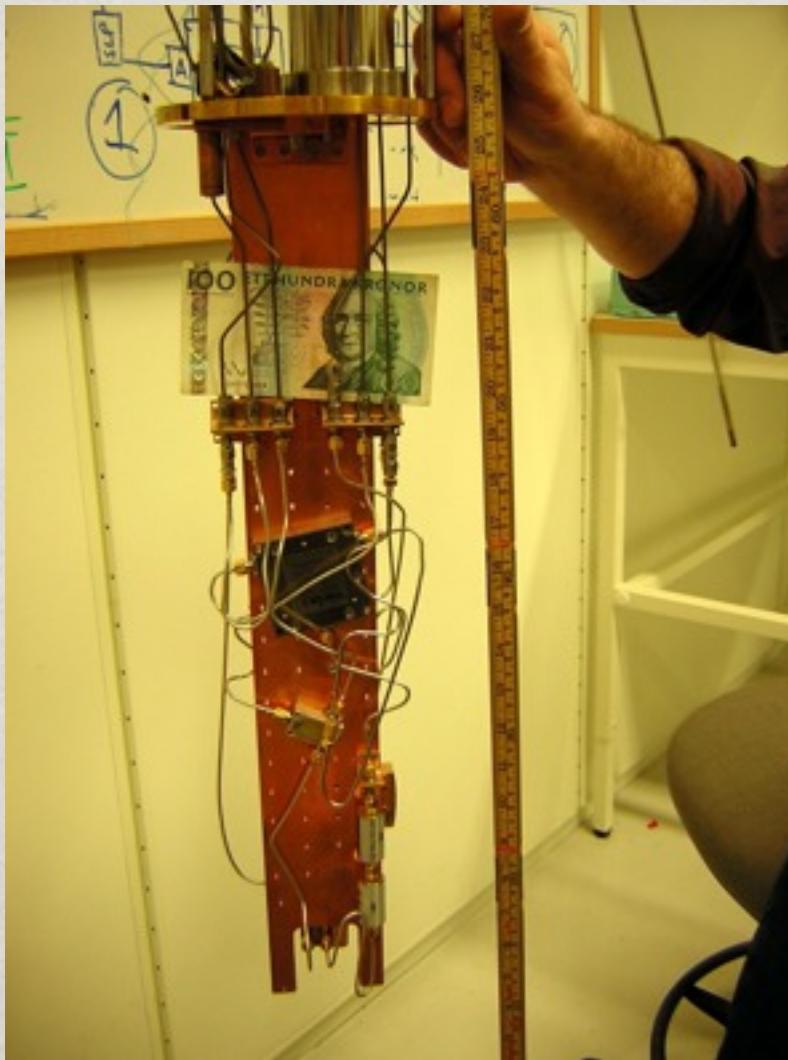
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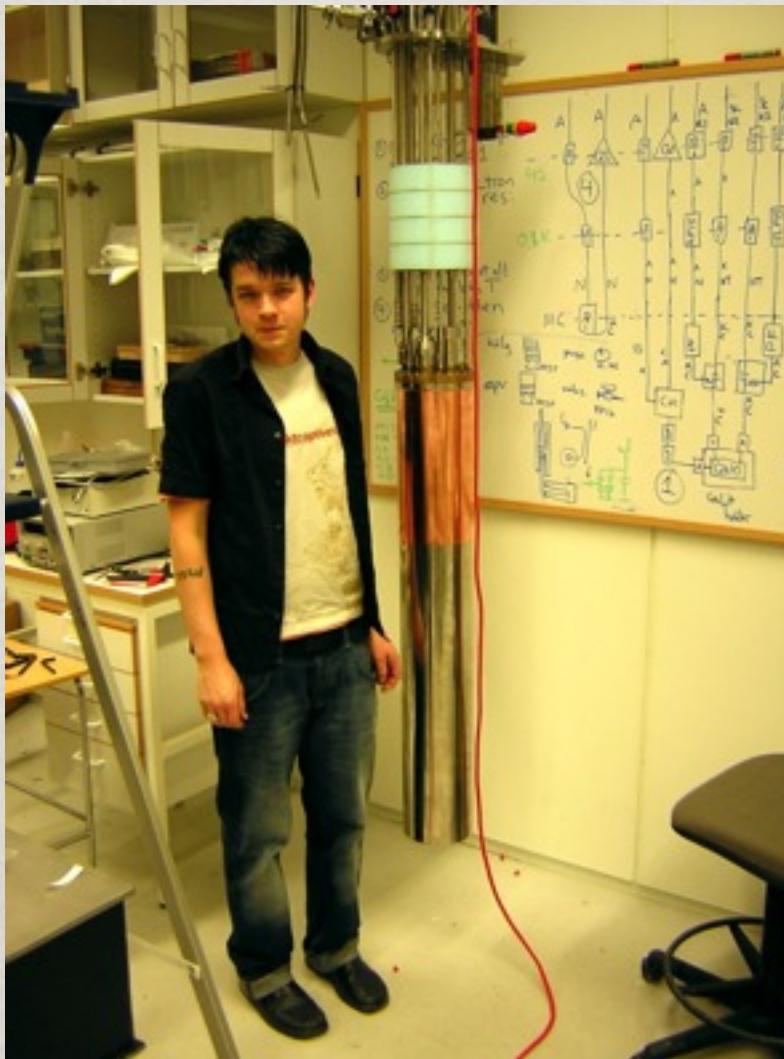
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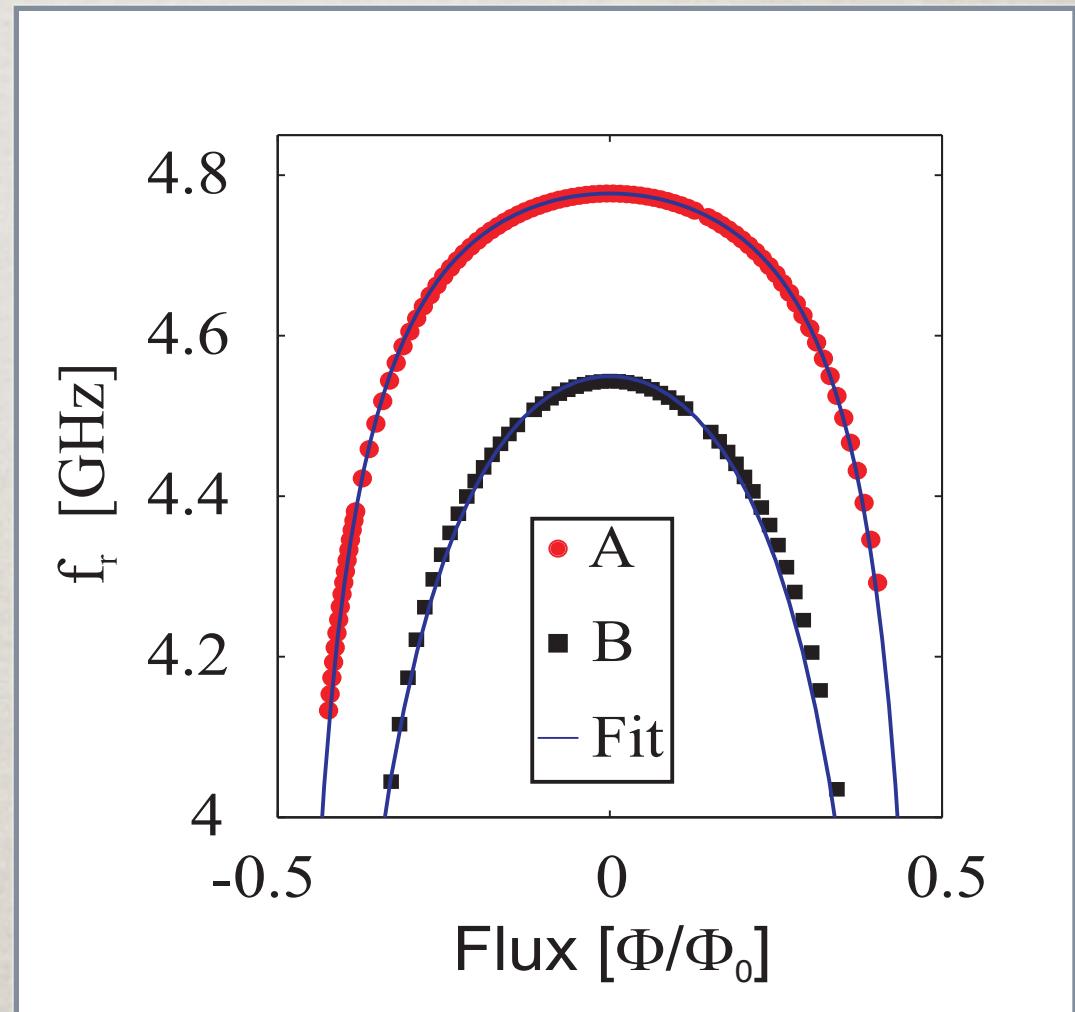
TUNABILITY

- Frequency tuning explained well by Josephson ind.

$$f(\Phi) = \frac{f_0}{1 + \frac{L_s(\Phi)}{Ll}}$$

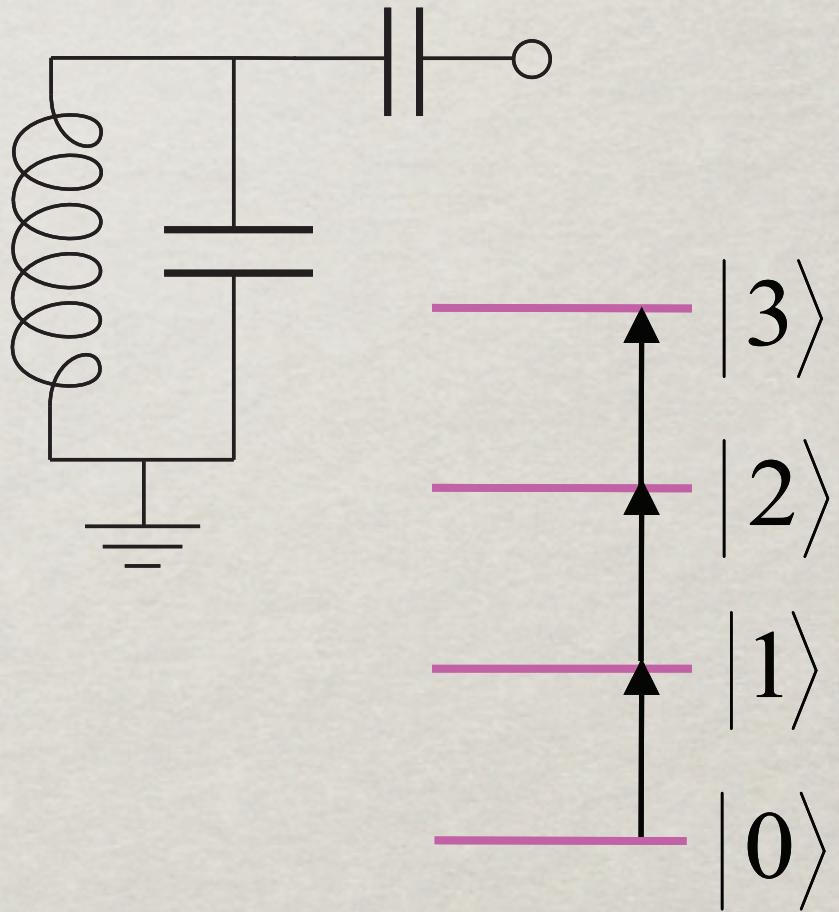
Sandberg, et al. *APL*, 92 (2008)

Yamamoto, et al. *APL*, 93 (2008)



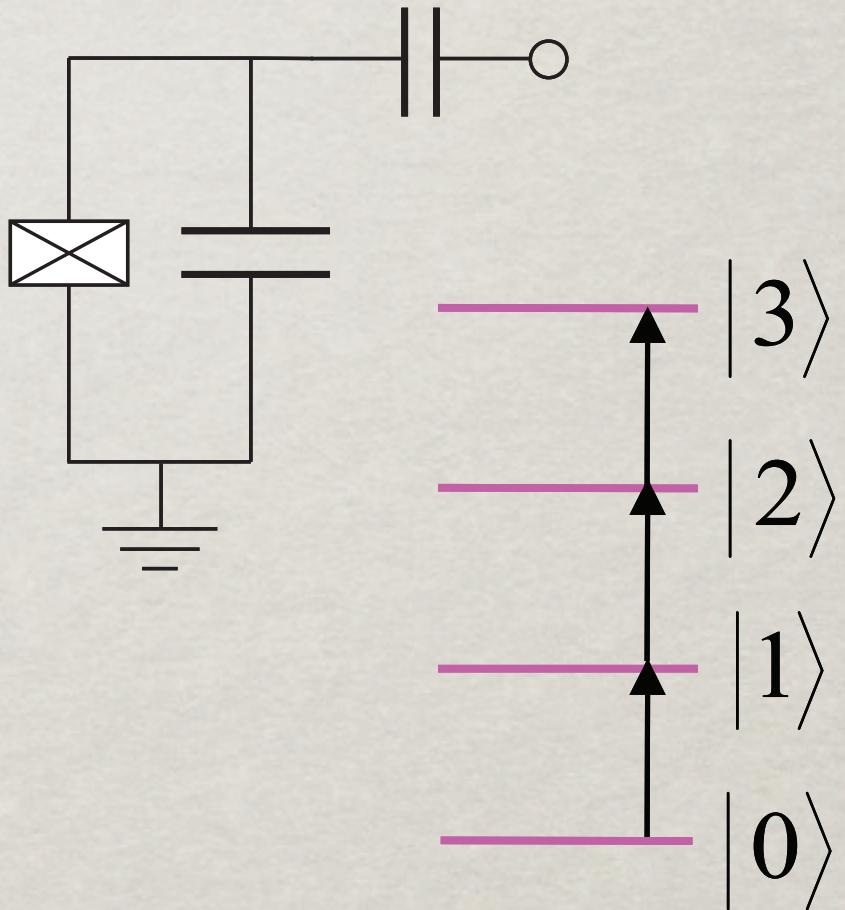
ARTIFICIAL ATOMS

- Quantized electrical circuit
- Harmonic oscillator is not a qubit
- Nonlinearity makes the circuit anharmonic and addressable
- Small JJ is a good nonlinear inductor



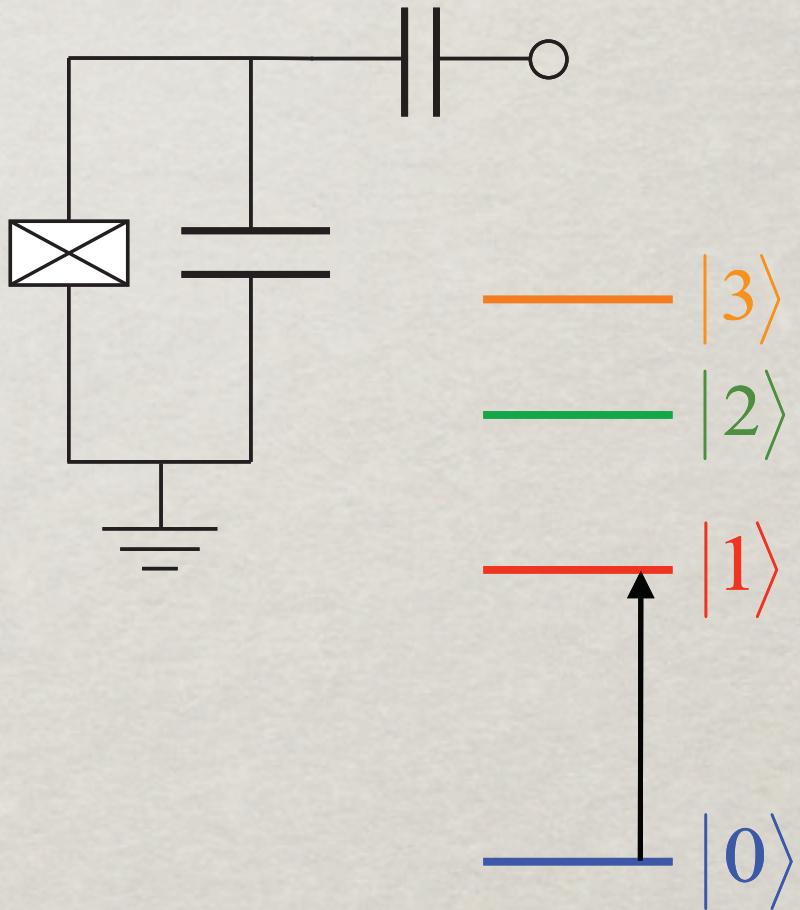
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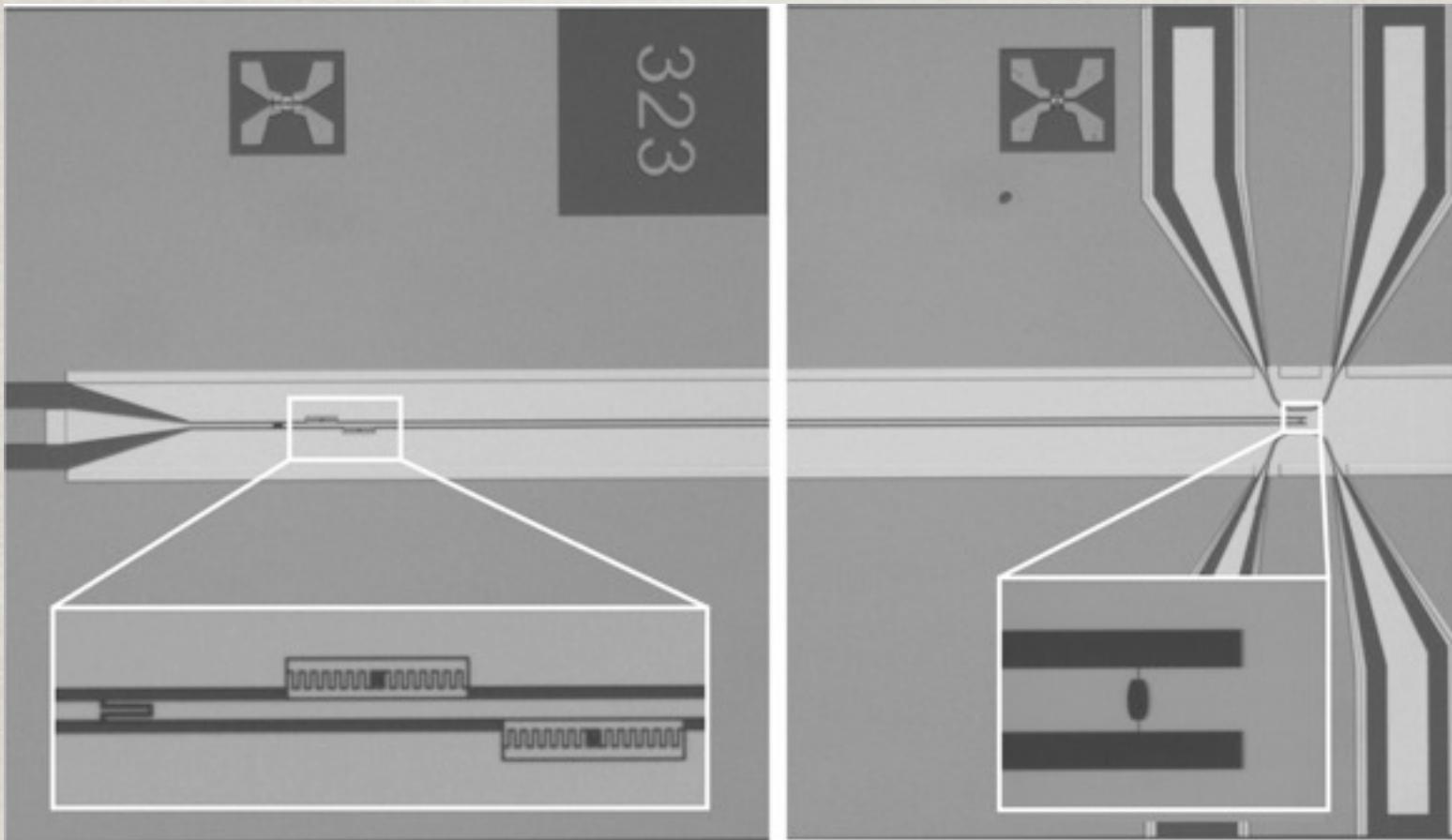


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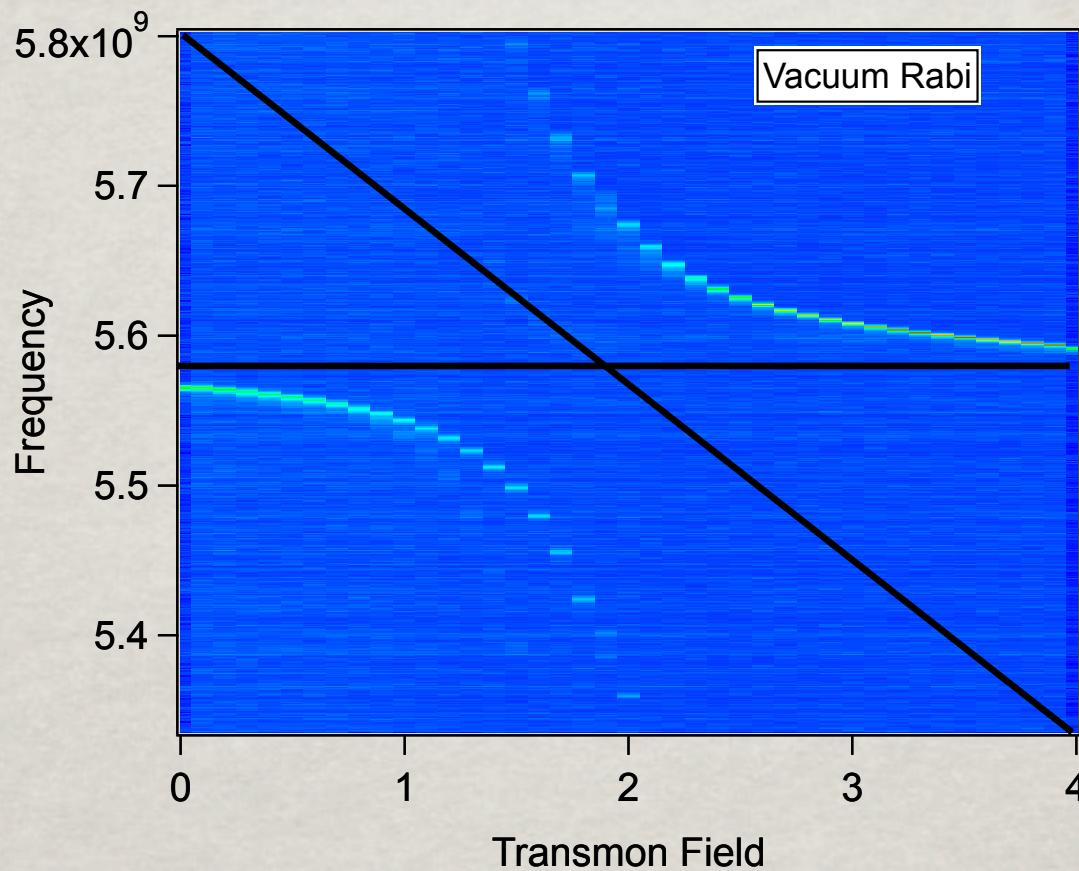


CIRCUIT QED



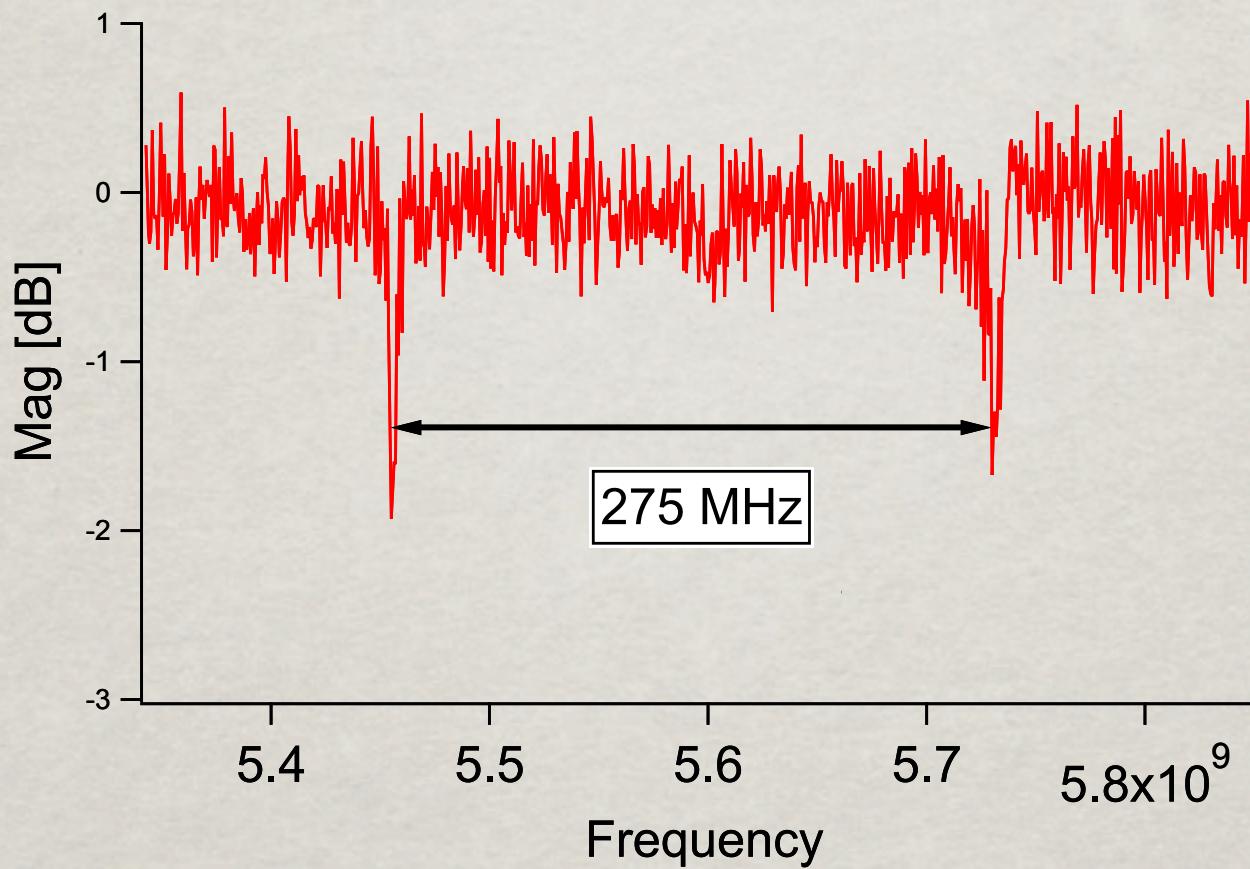
Wallraff et al, Nature 431 (2004) (and many more)
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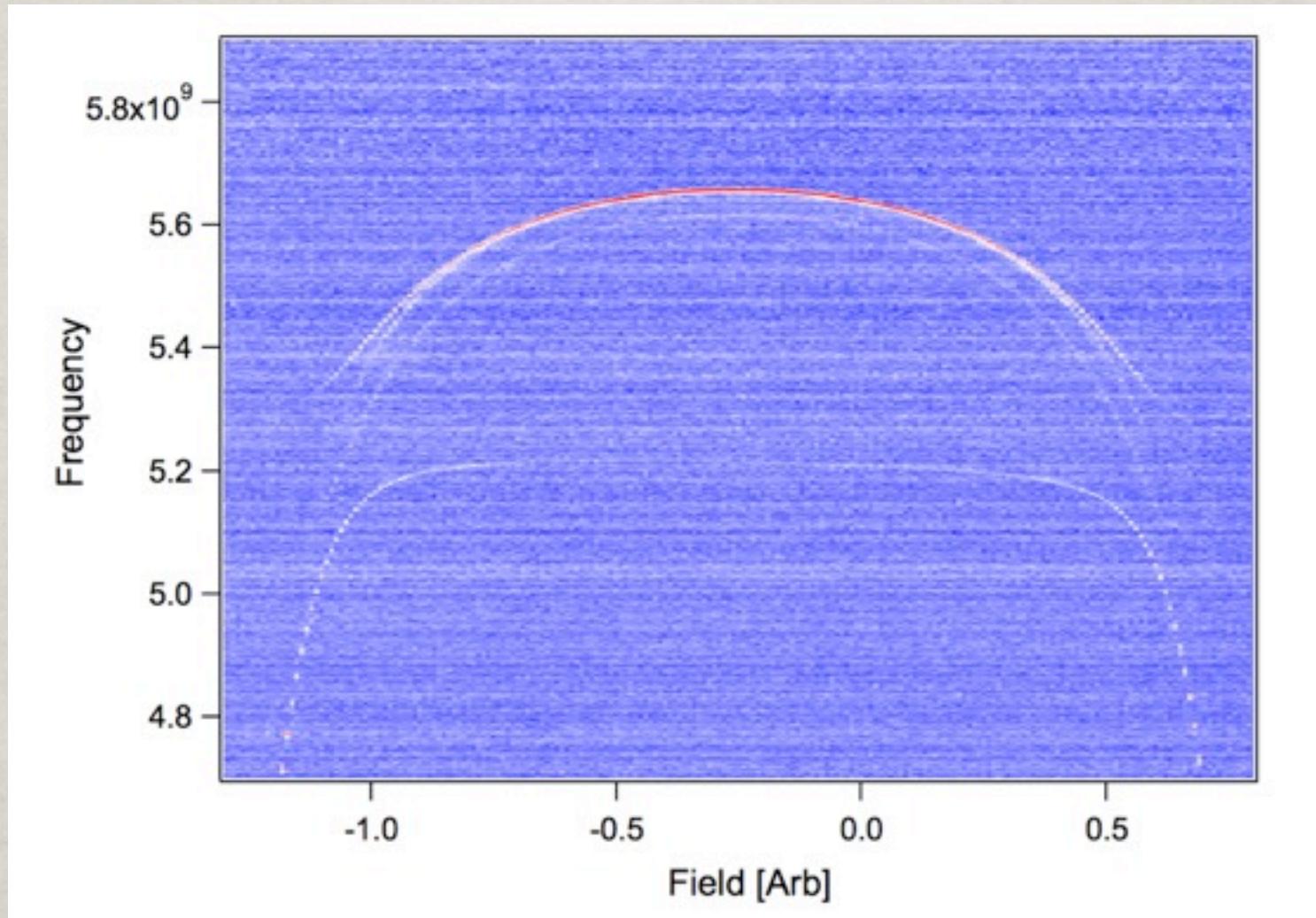
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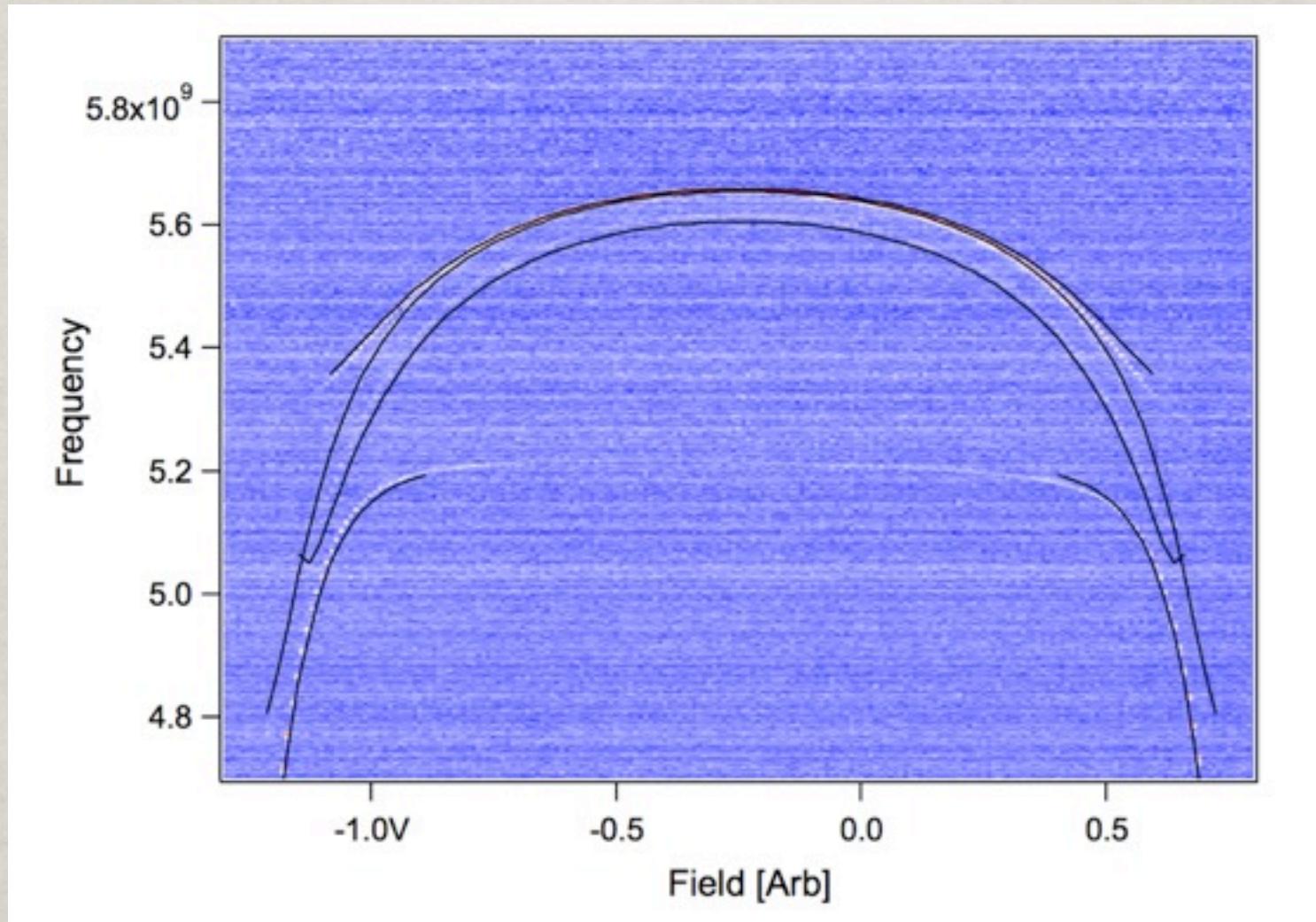
VACUUM RABI

✿ Tuning the resonator

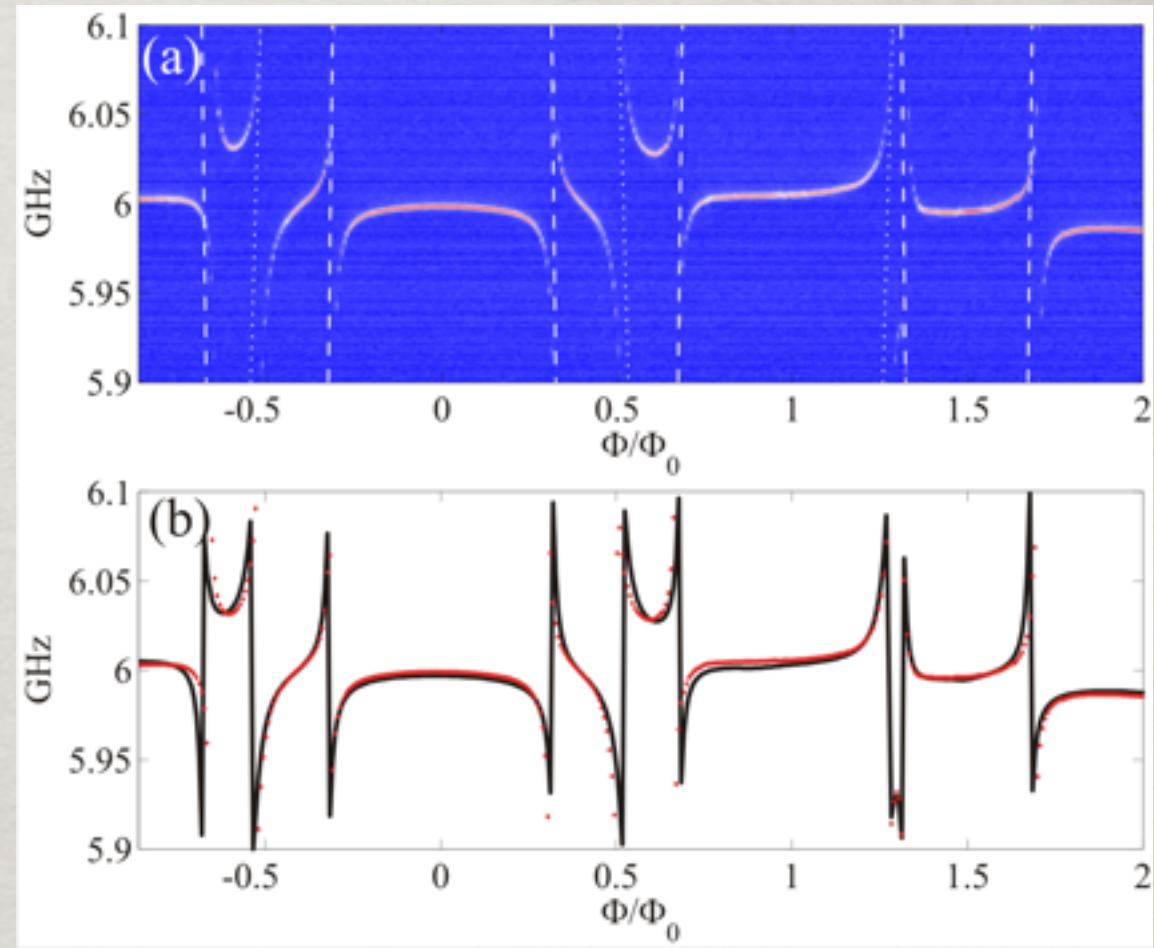
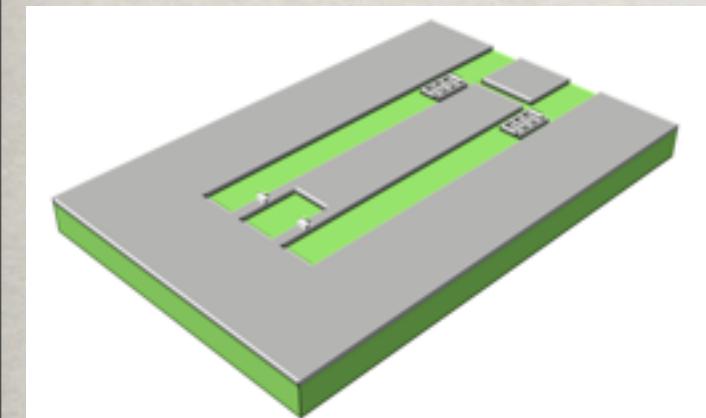


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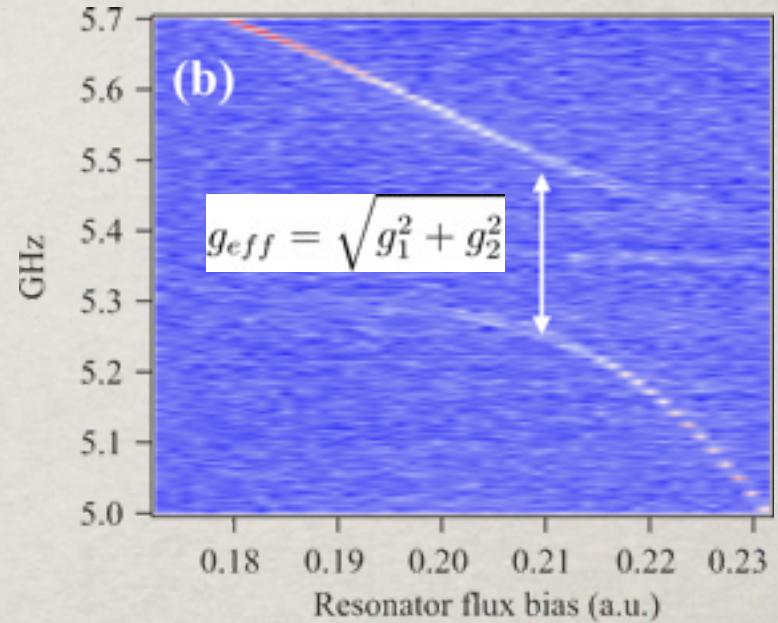
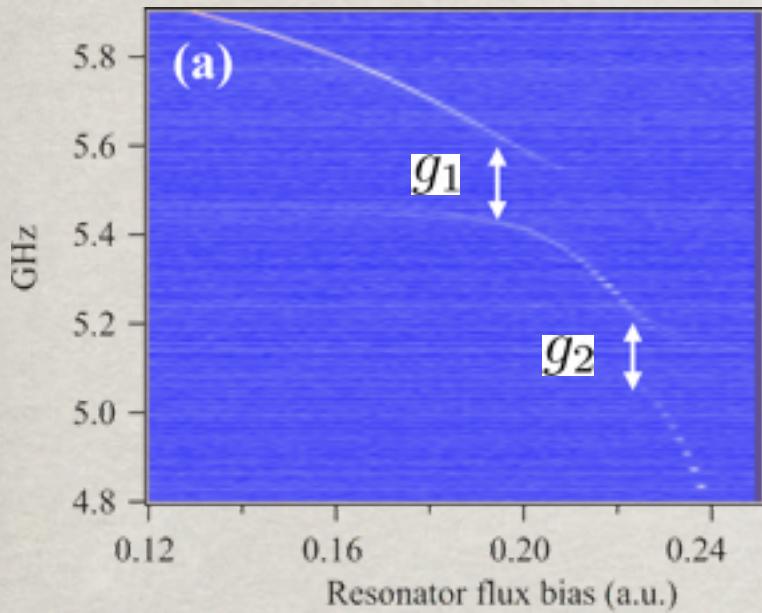


2-QUBIT CIRCUIT QED



- ✿ Fit only the flux coupling, others parameters from spectroscopy

2-QUBIT CIRCUIT QED



- ✿ Tune resonator with qubits detuned

- ✿ Now with qubits on resonance
- ✿ Qubits acts as one collective system: Tavis-Cummings model

SINGLE-ATOM SCATTERING AND 1D QED

SINGLE-ATOM SCATTERING

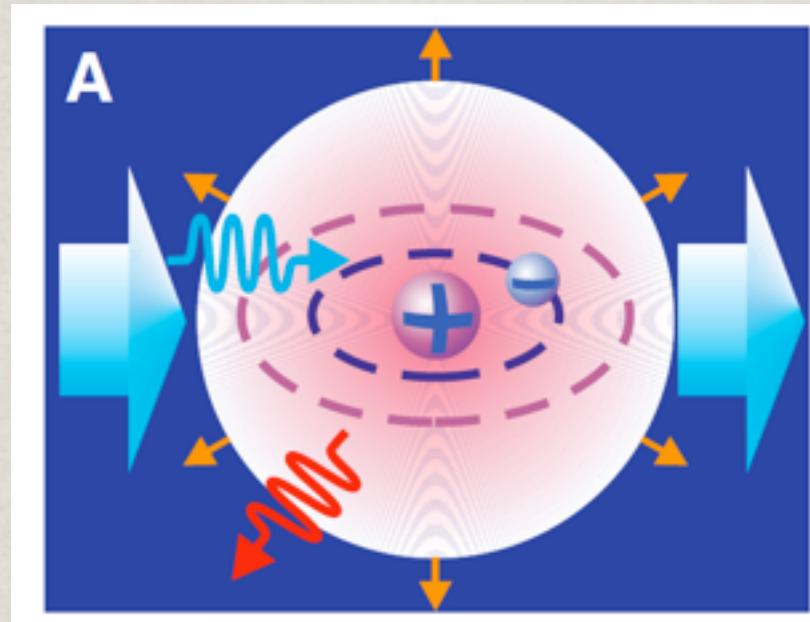


Fig: Astafiev et al., Science (2010).

SINGLE-ATOM SCATTERING

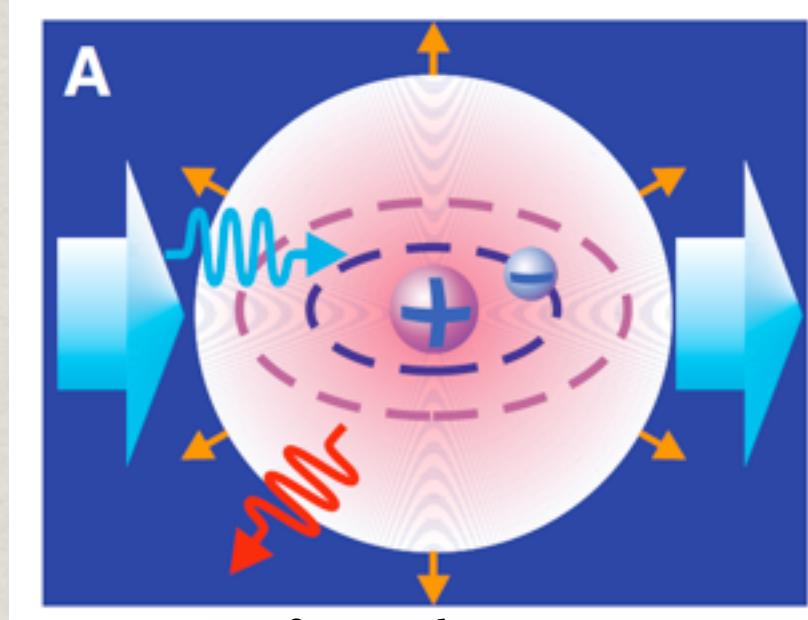


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- ✿ What is the maximum reflection from a single atom?

SINGLE-ATOM SCATTERING

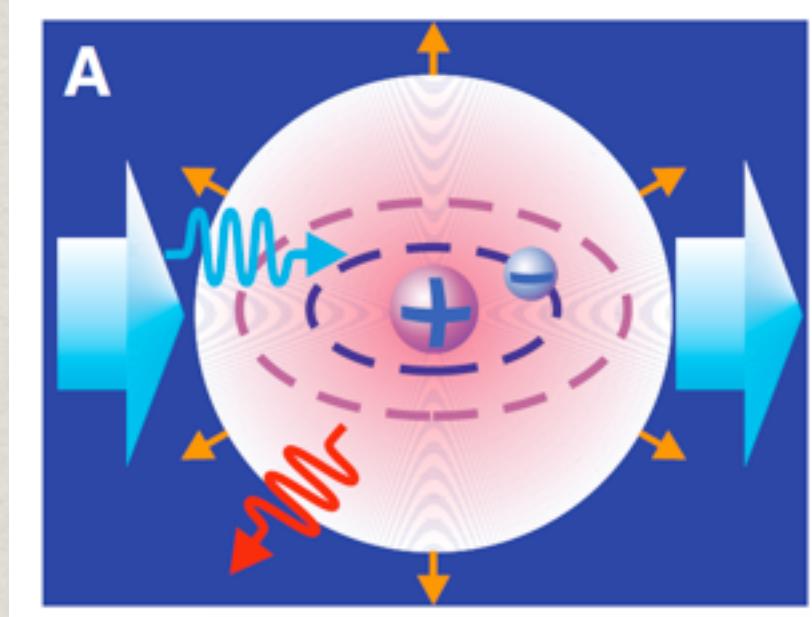


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 - ✿ First guess: 50% because 1/2 of spontaneous emission goes forward, 1/2 back

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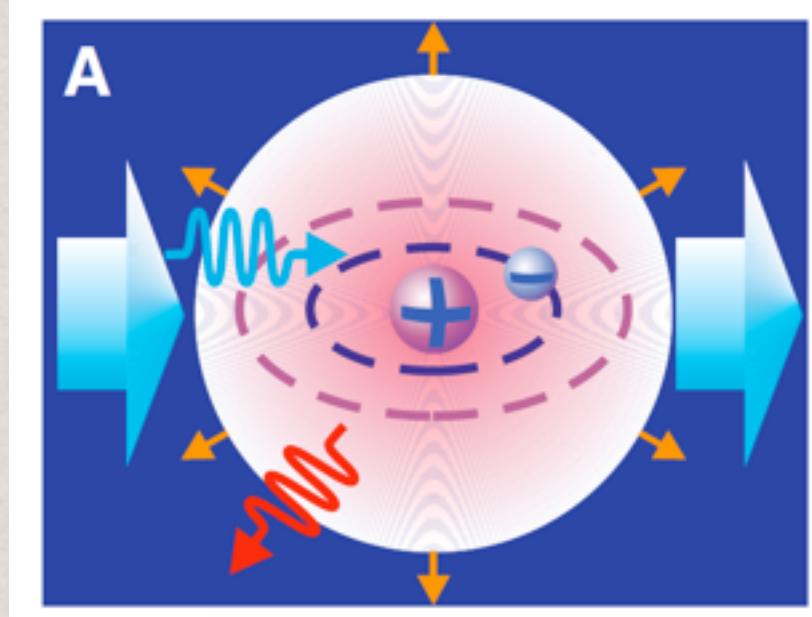
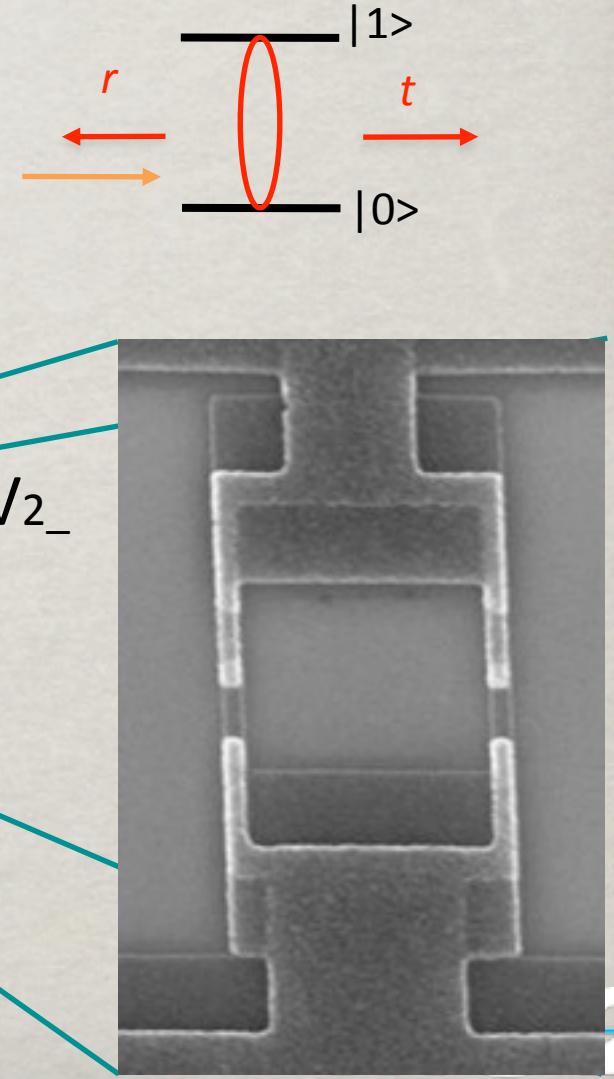
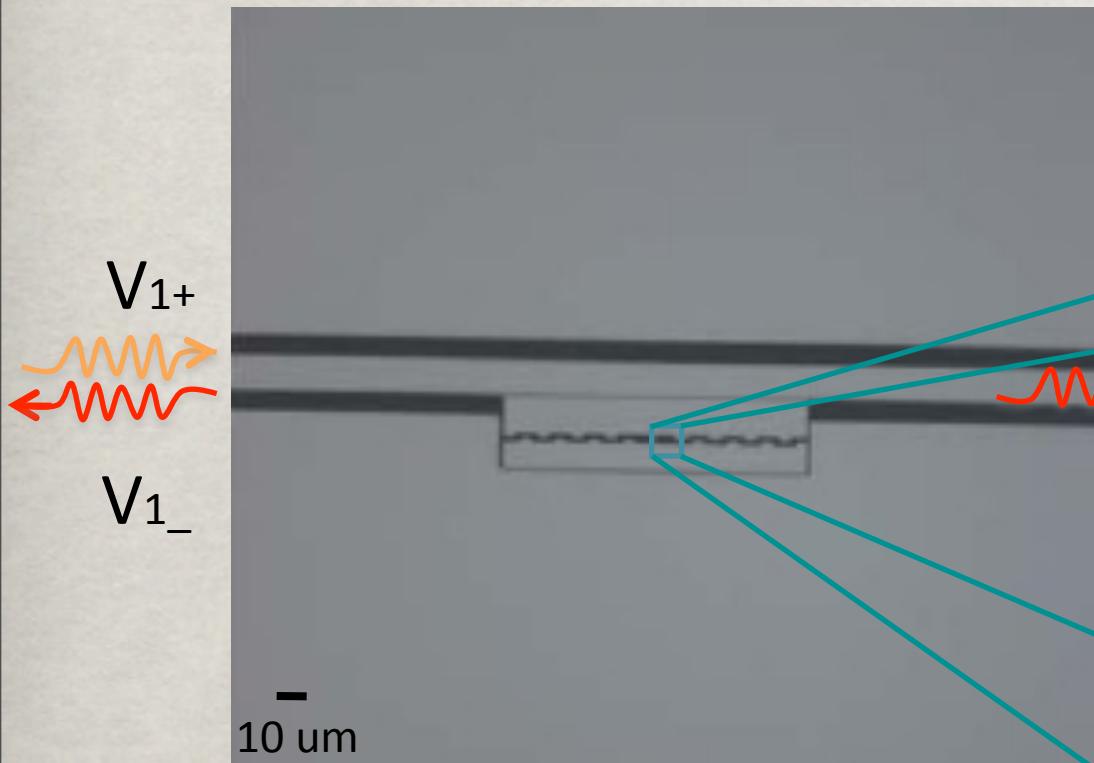


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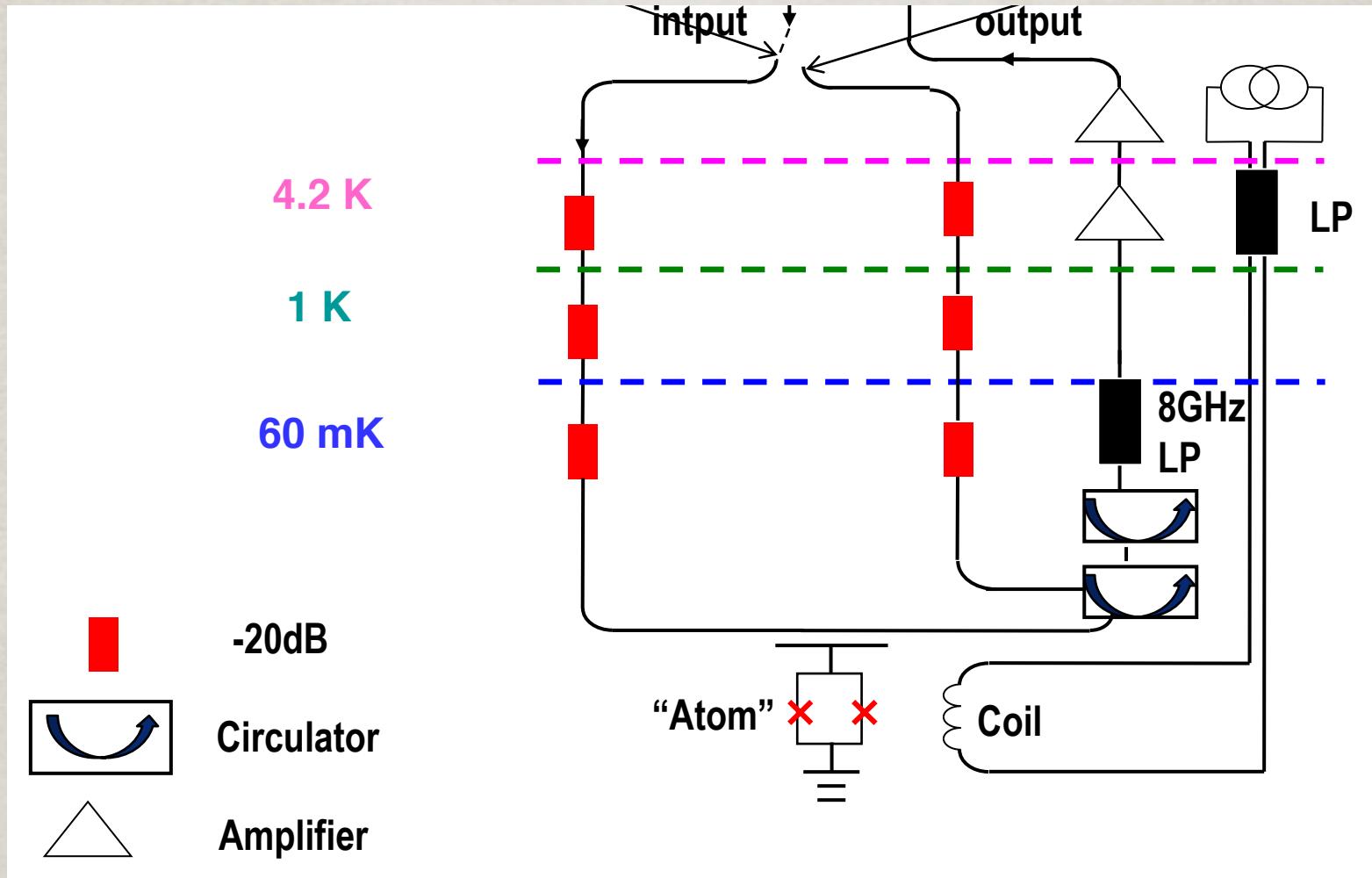
- ✿ What is the maximum reflection from a single atom?
 - ✿ First guess: 50% because 1/2 of spontaneous emission goes forward, 1/2 back
 - ✿ Fully coherent: 100% because forward scattering destructively interferes with incoming wave

ATOM IN A TRANSMISSION LINE

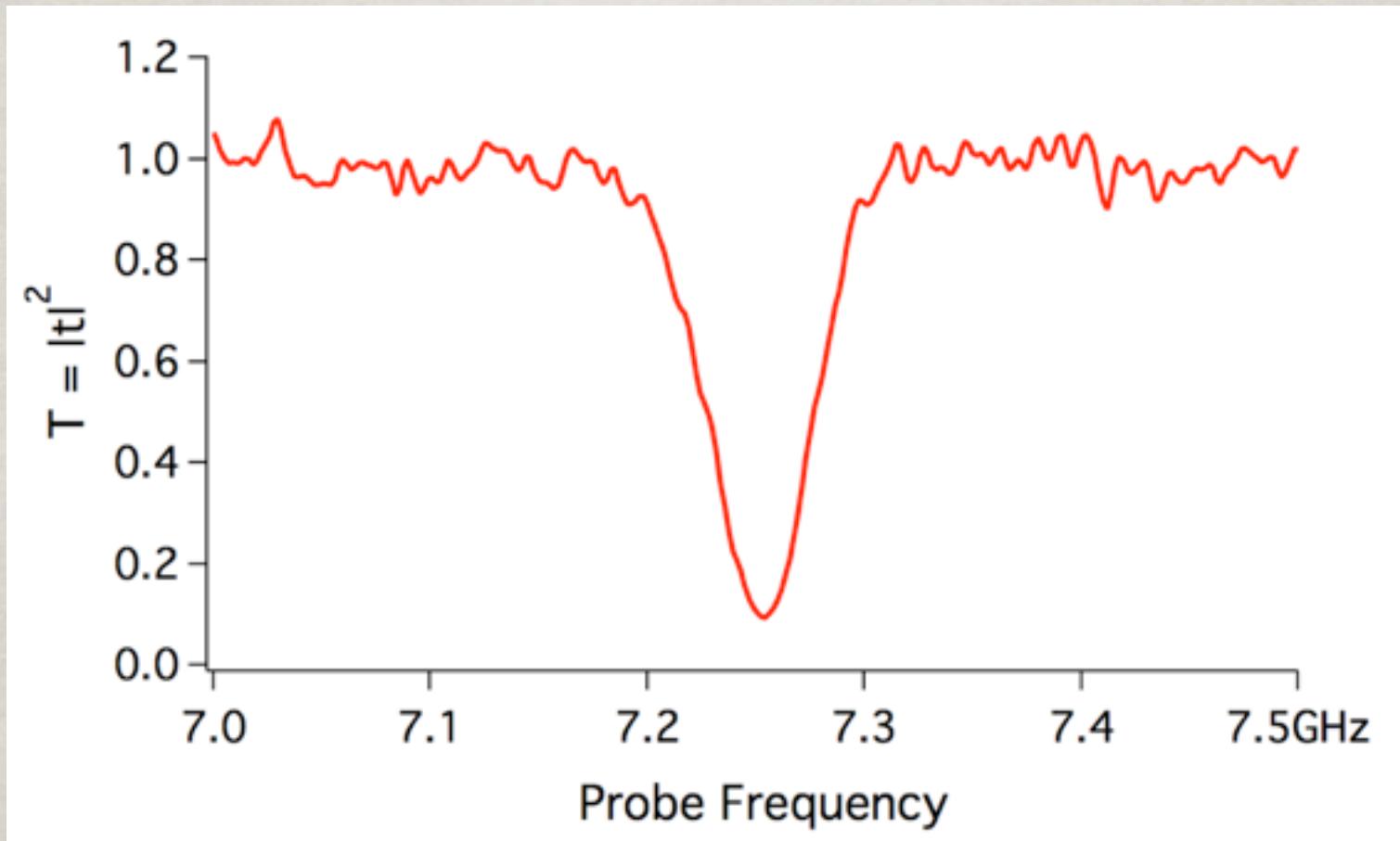
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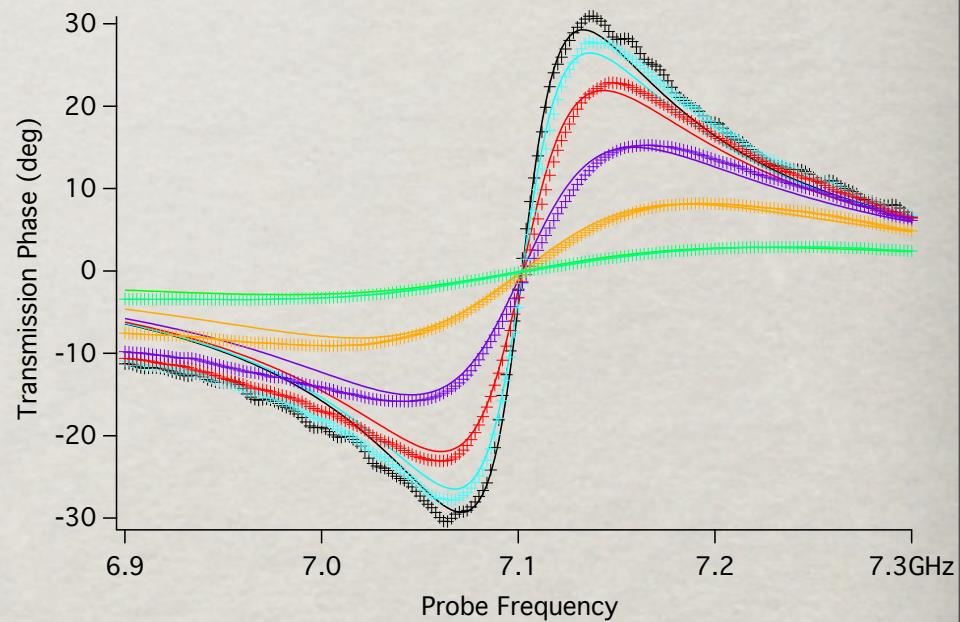
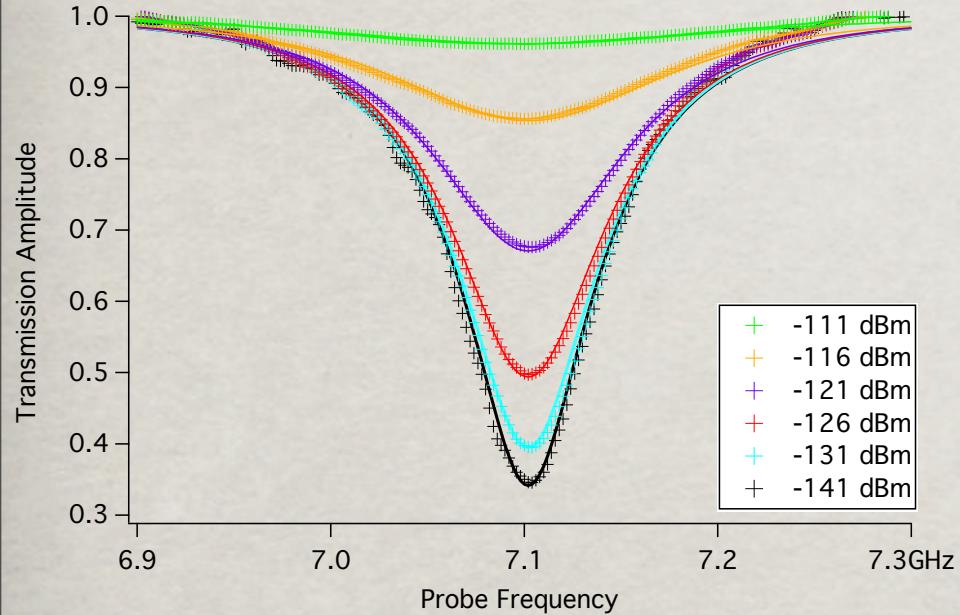
MEASUREMENT SETUP



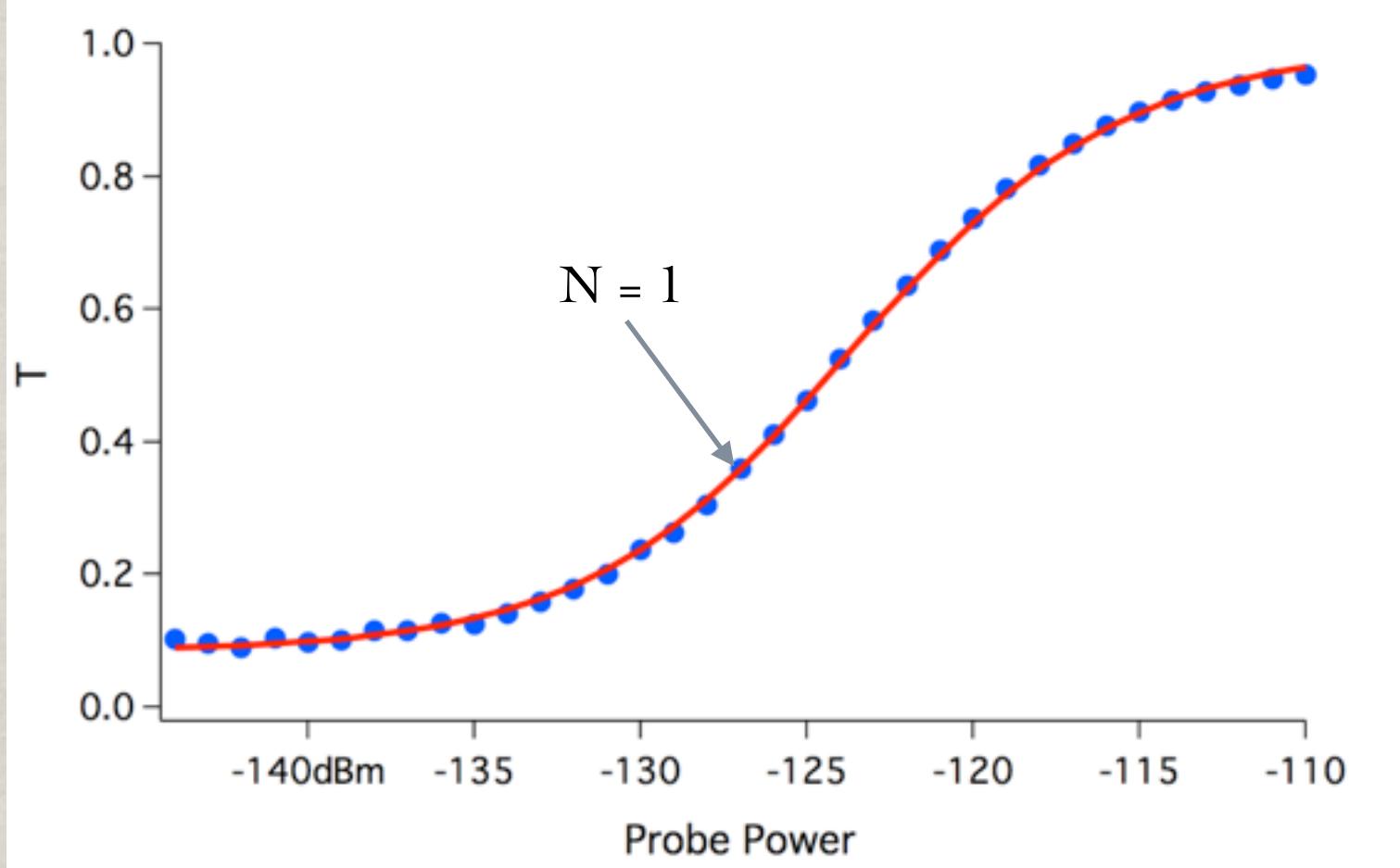
STRONG EXTINCTION



STRONG SATURATION

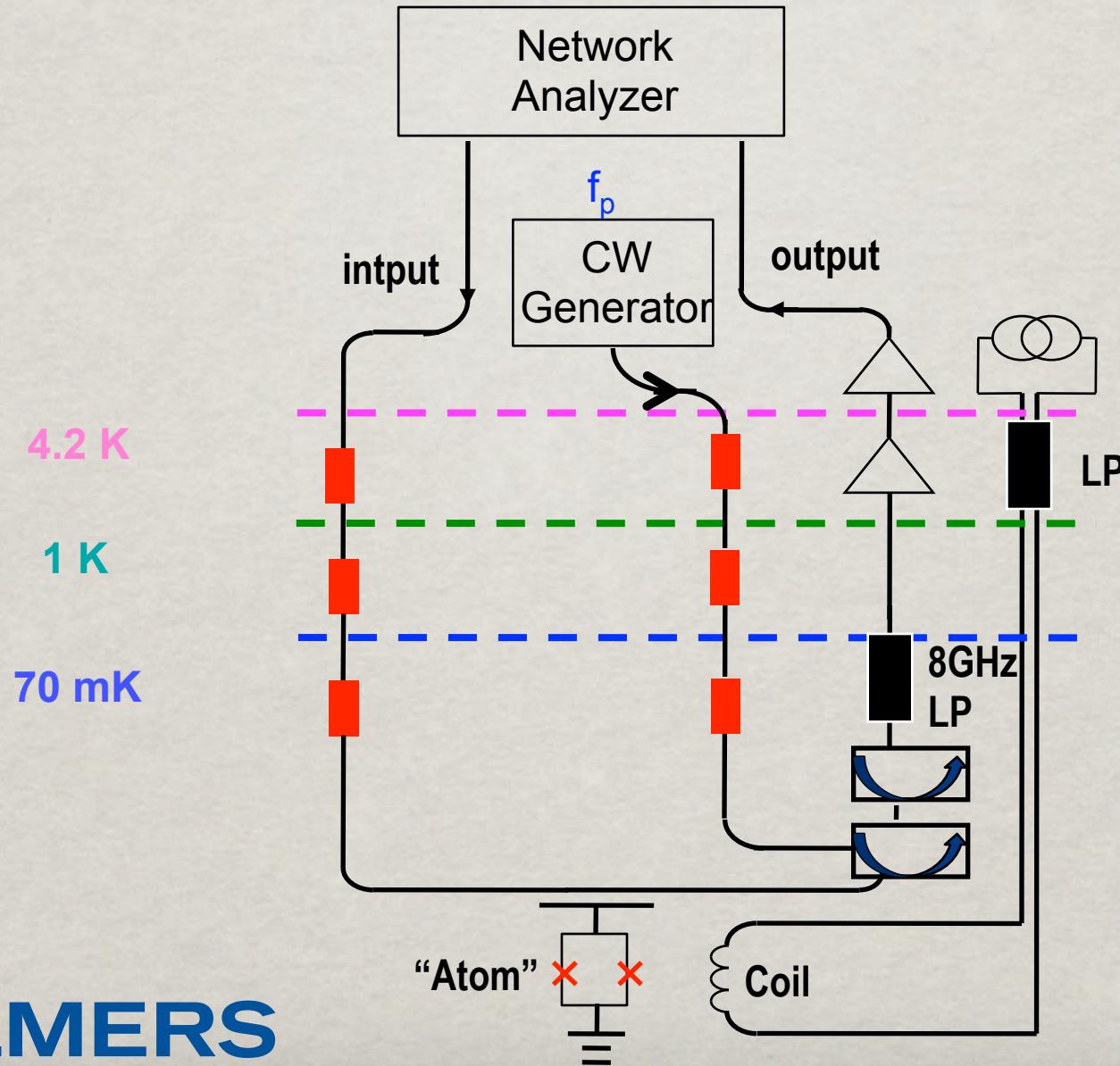


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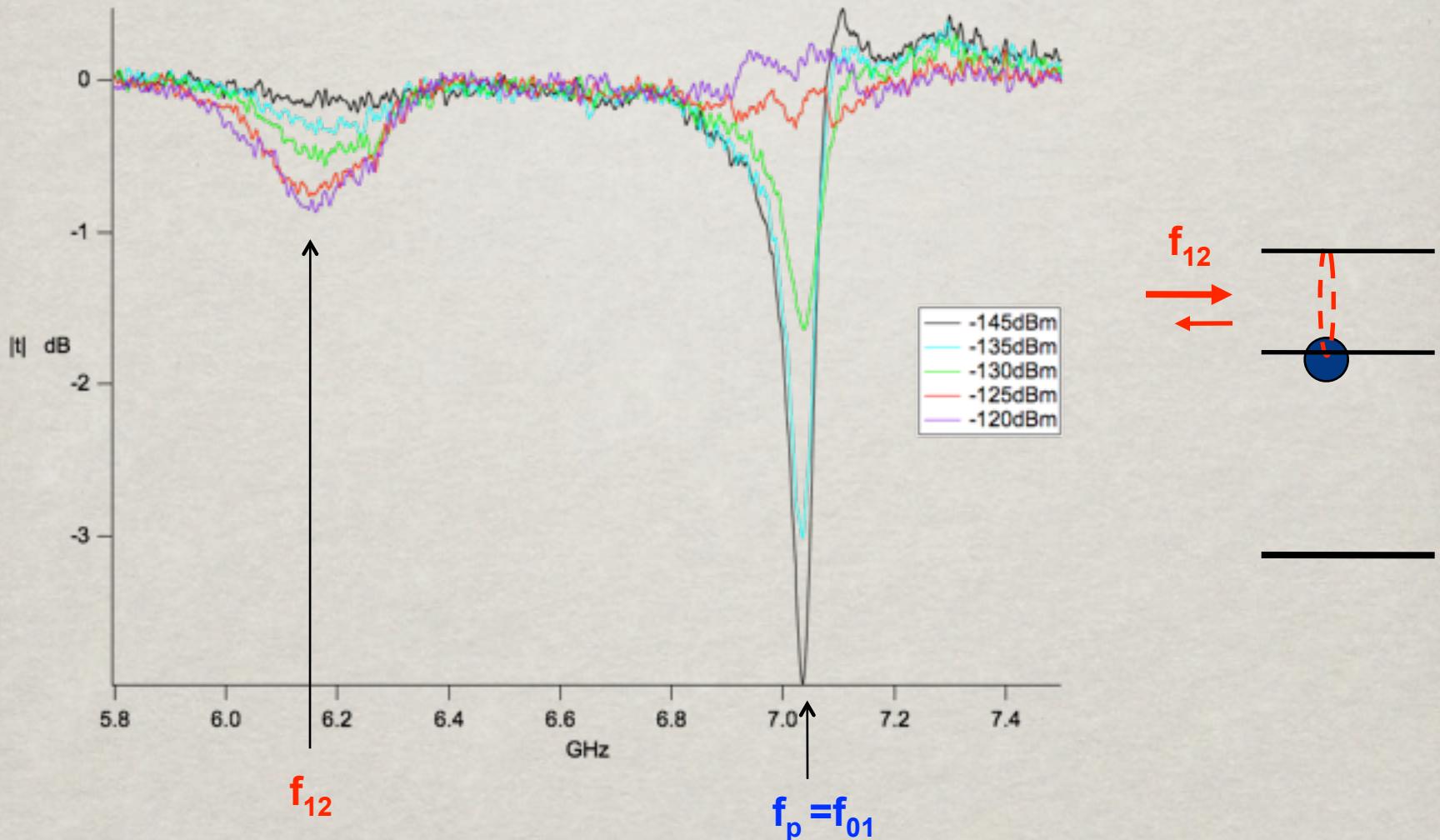


$$N = \frac{P}{hf_0\Gamma_1}$$

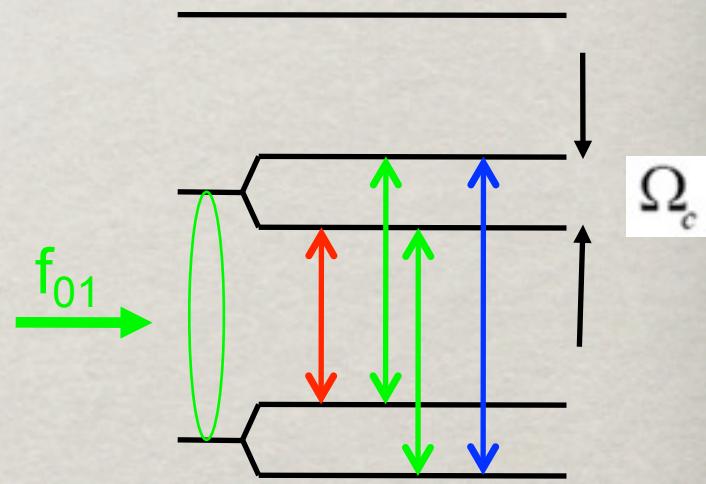
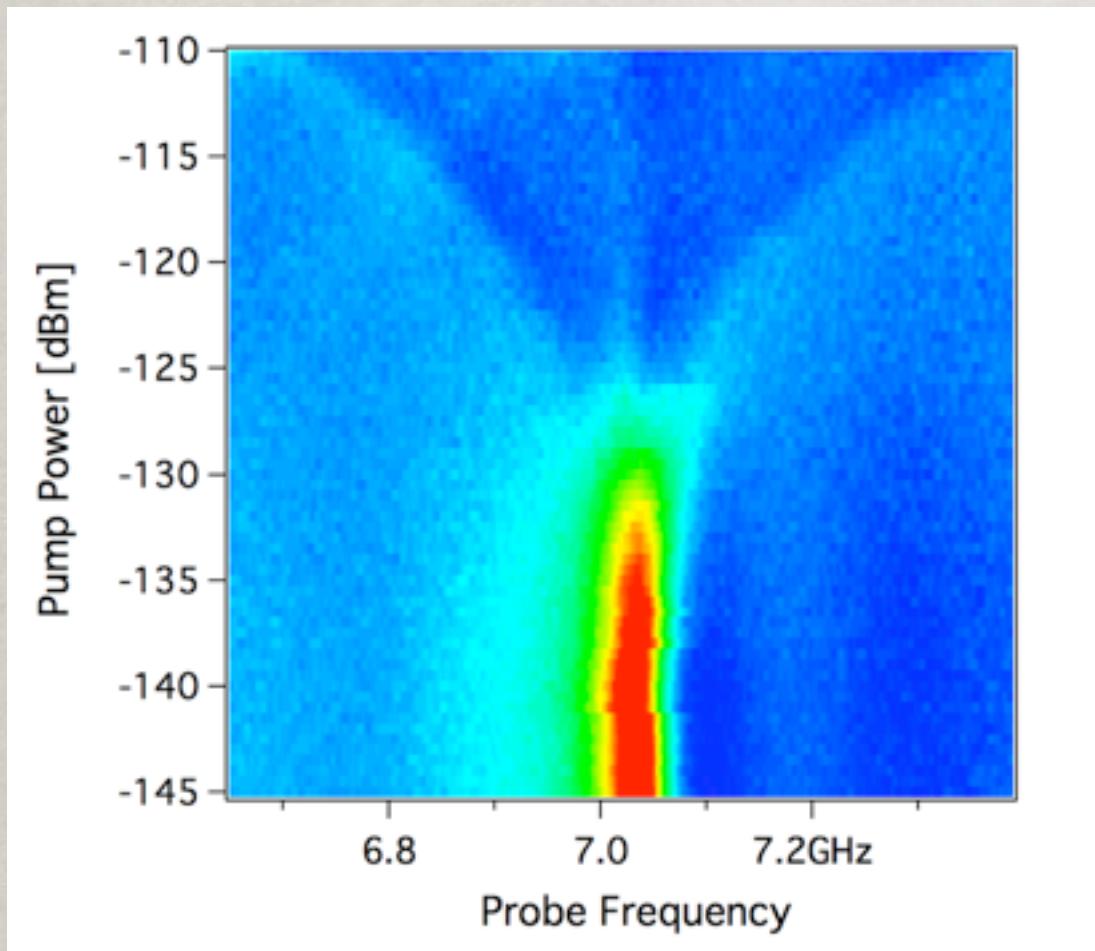
TWO-TONE SPECTROSCOPY



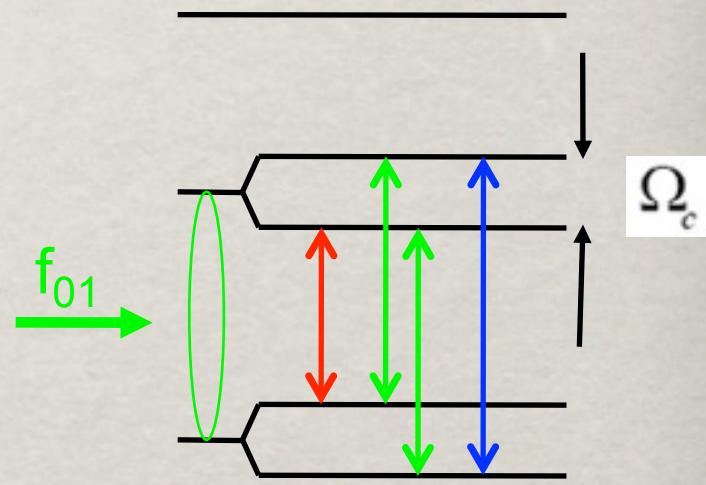
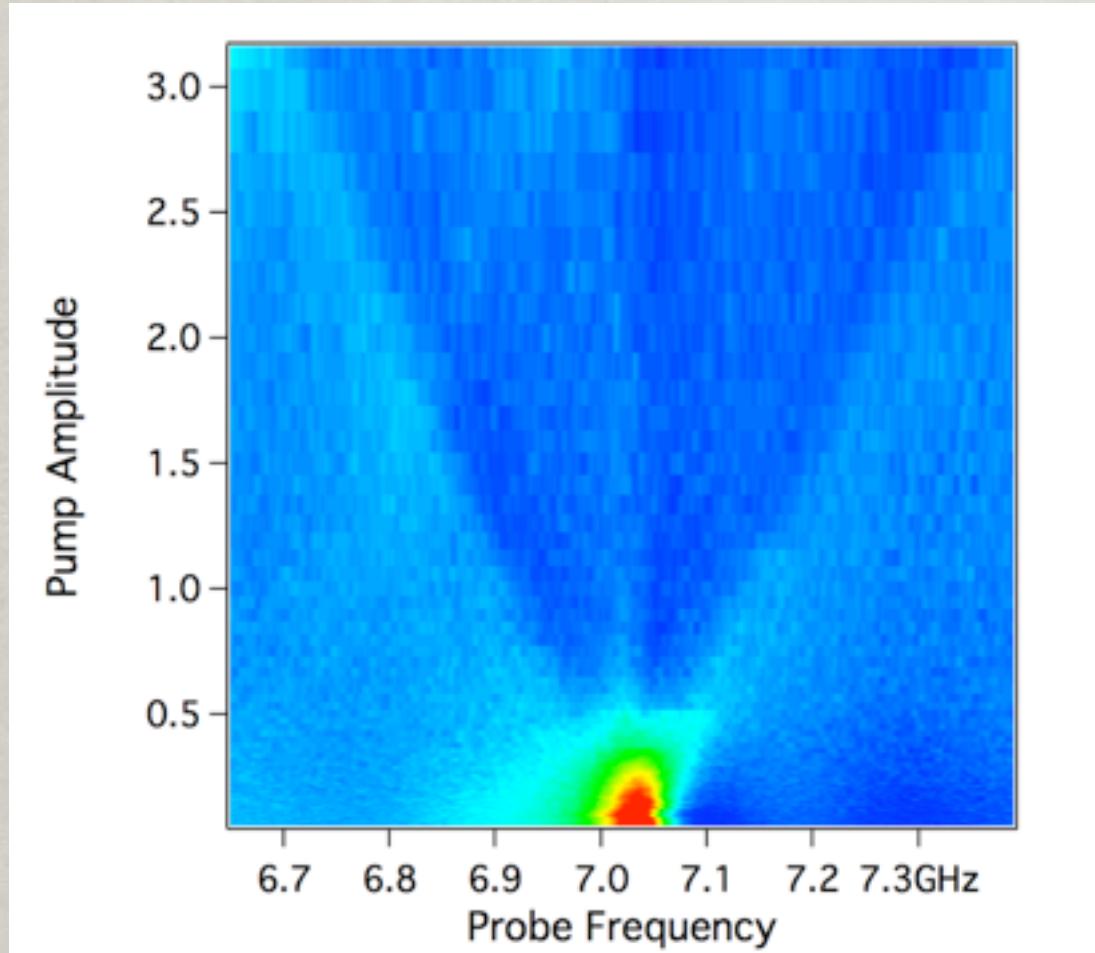
EXCITED EXTINCTION



MOLLOW TRIPLET

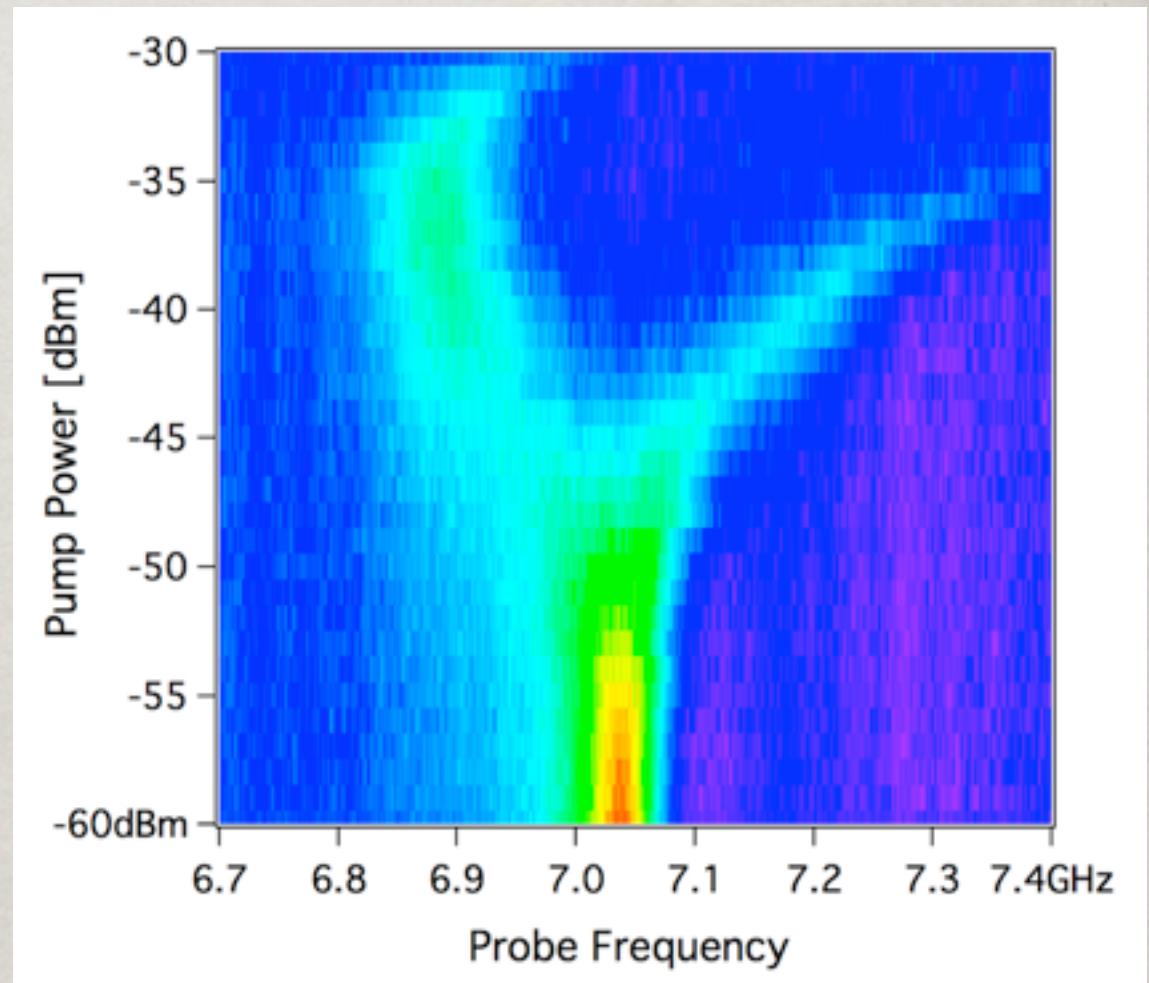
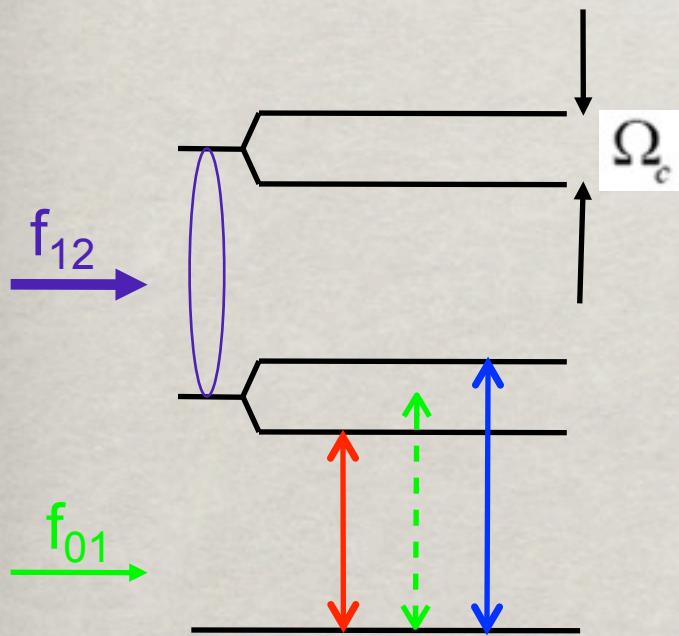


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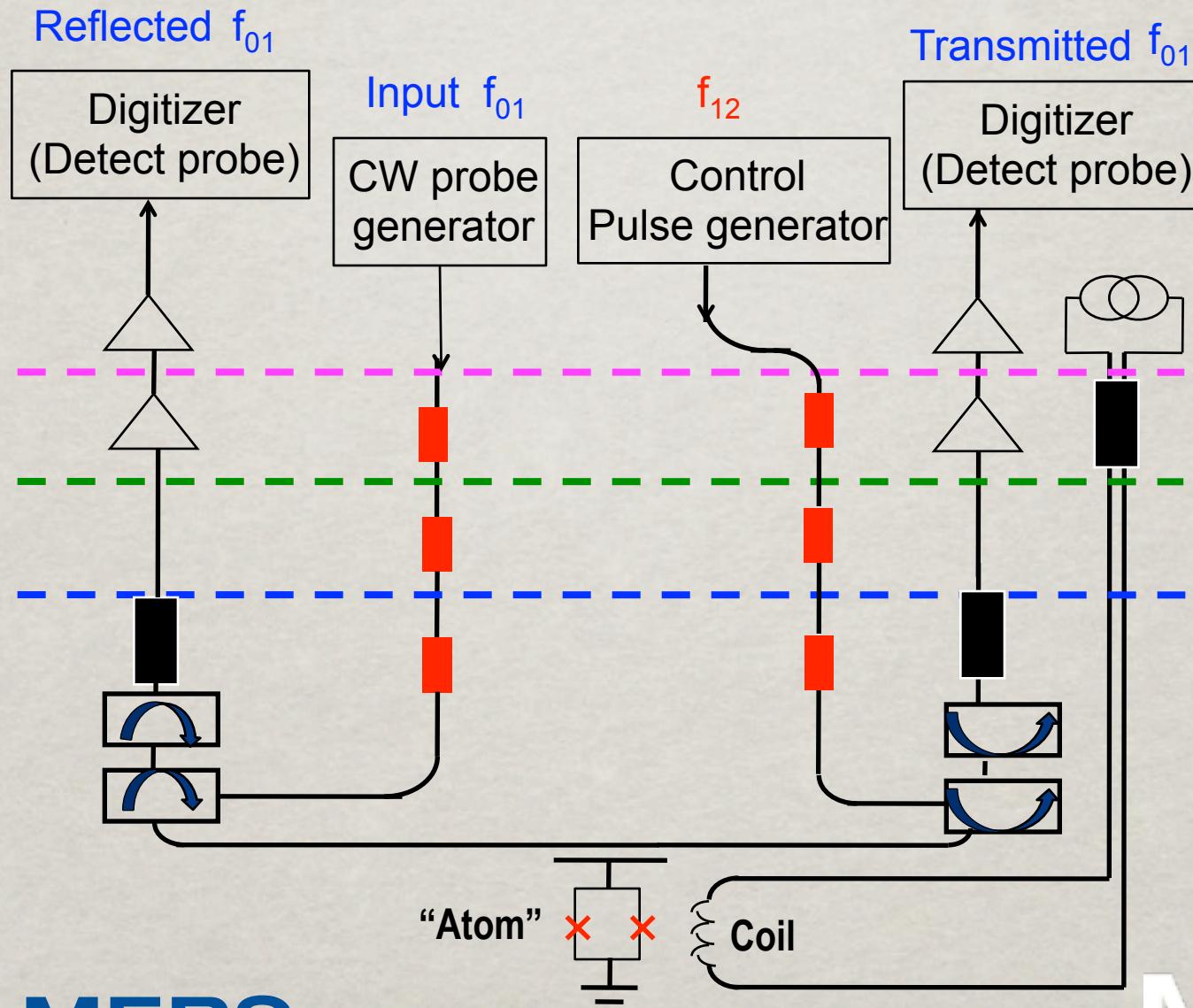


ELECTROMAGNETICALLY INDUCED TRANSPARENCY

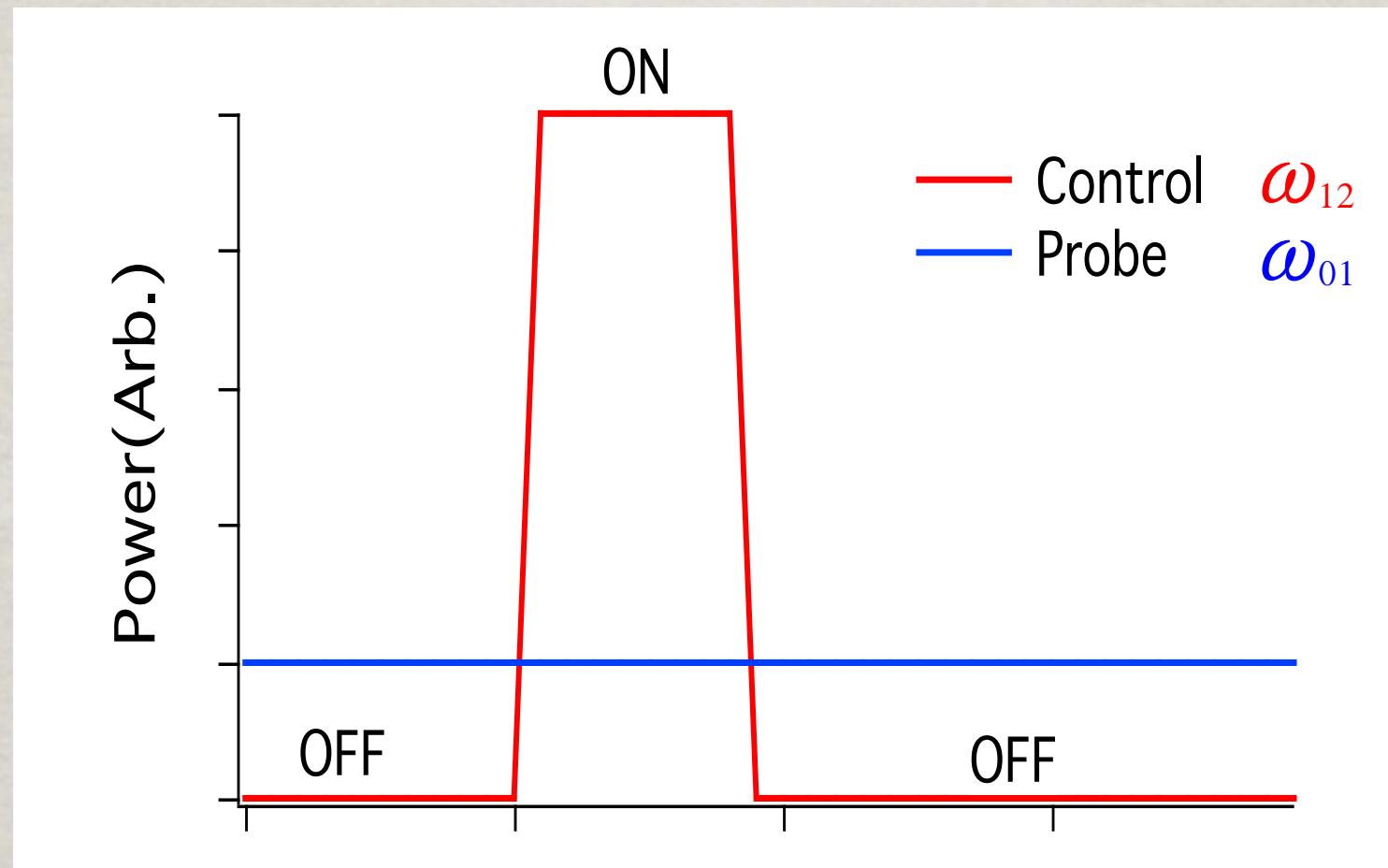
- ✿ Pump 1-2 transition



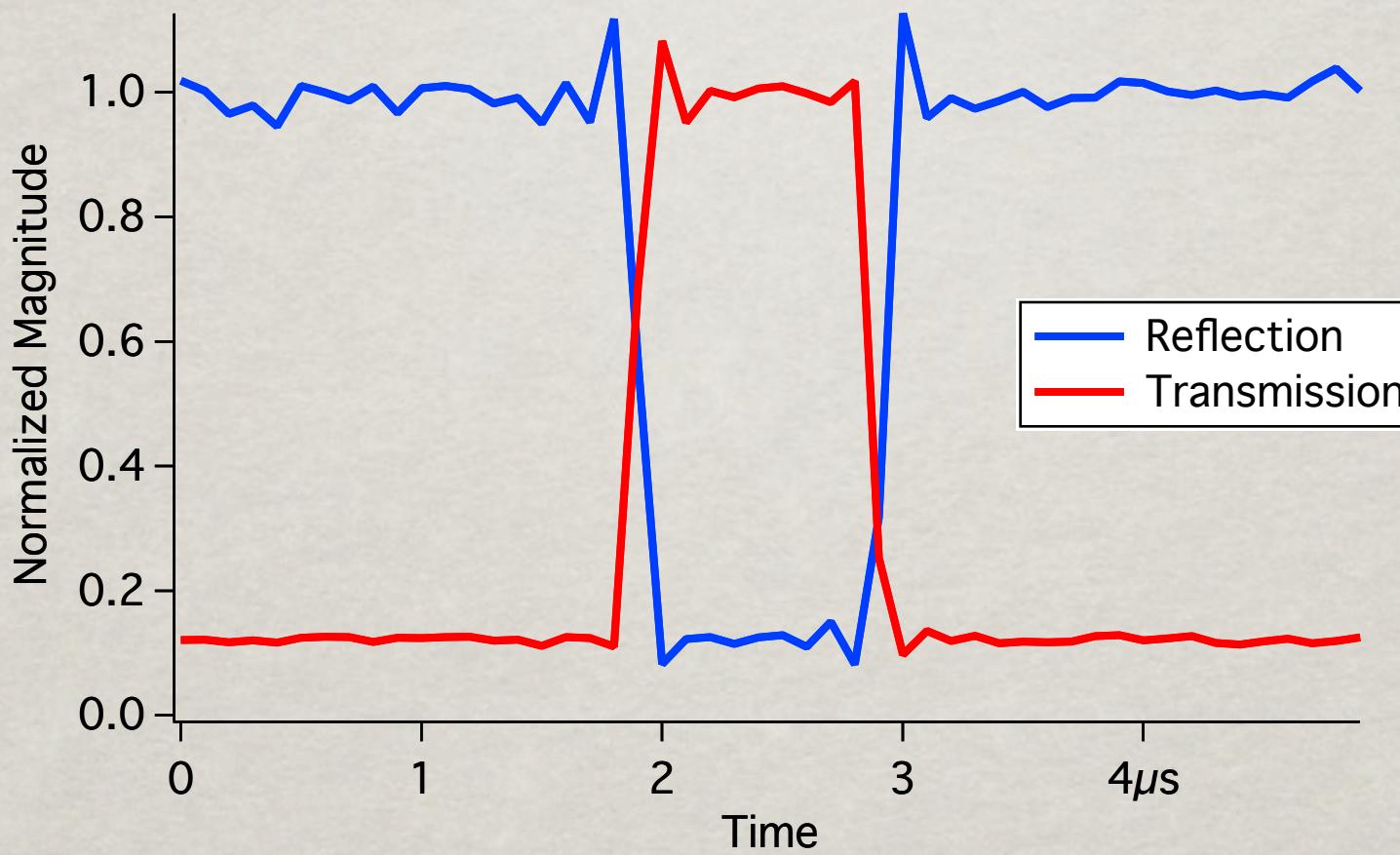
SINGLE-PHOTON ROUTER



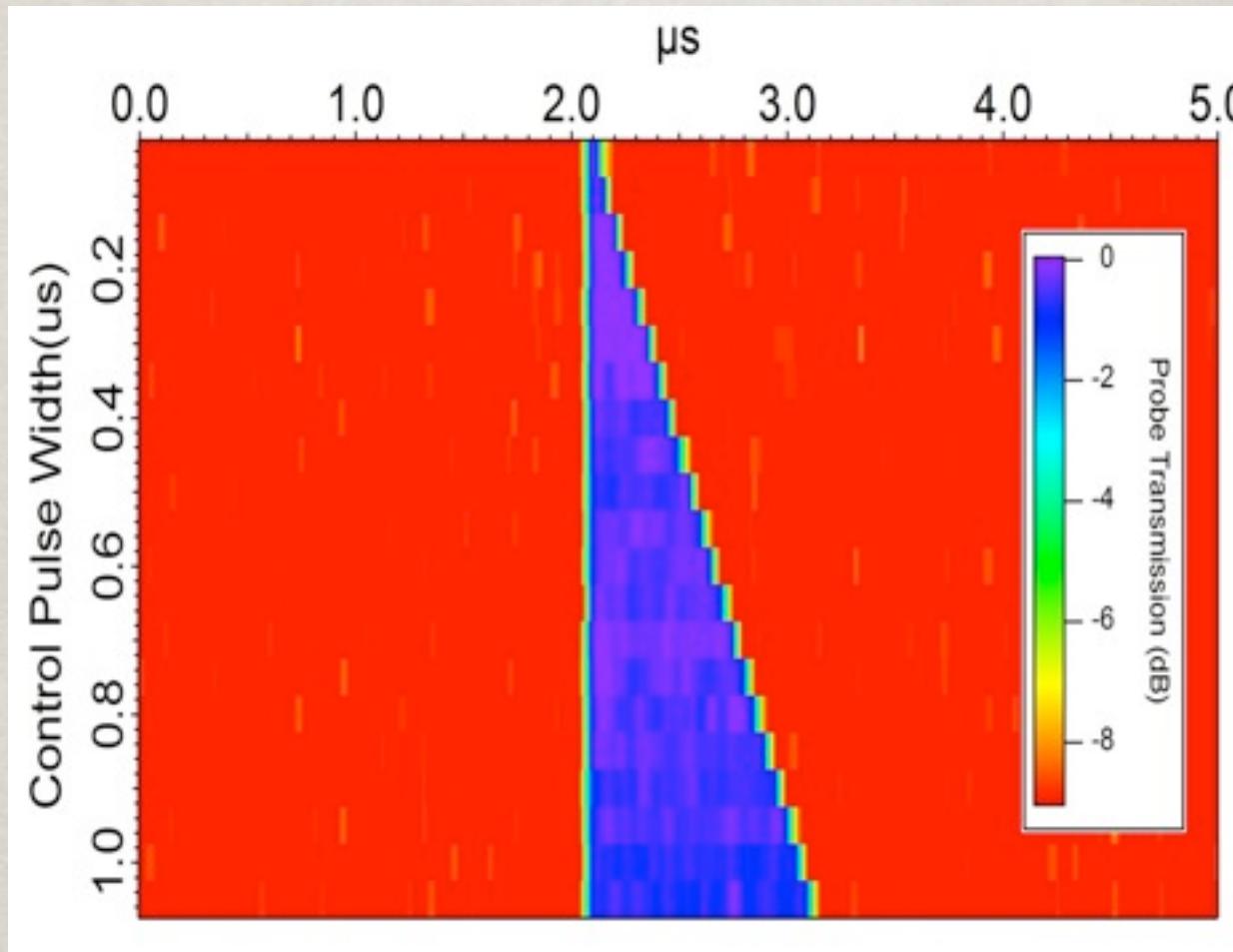
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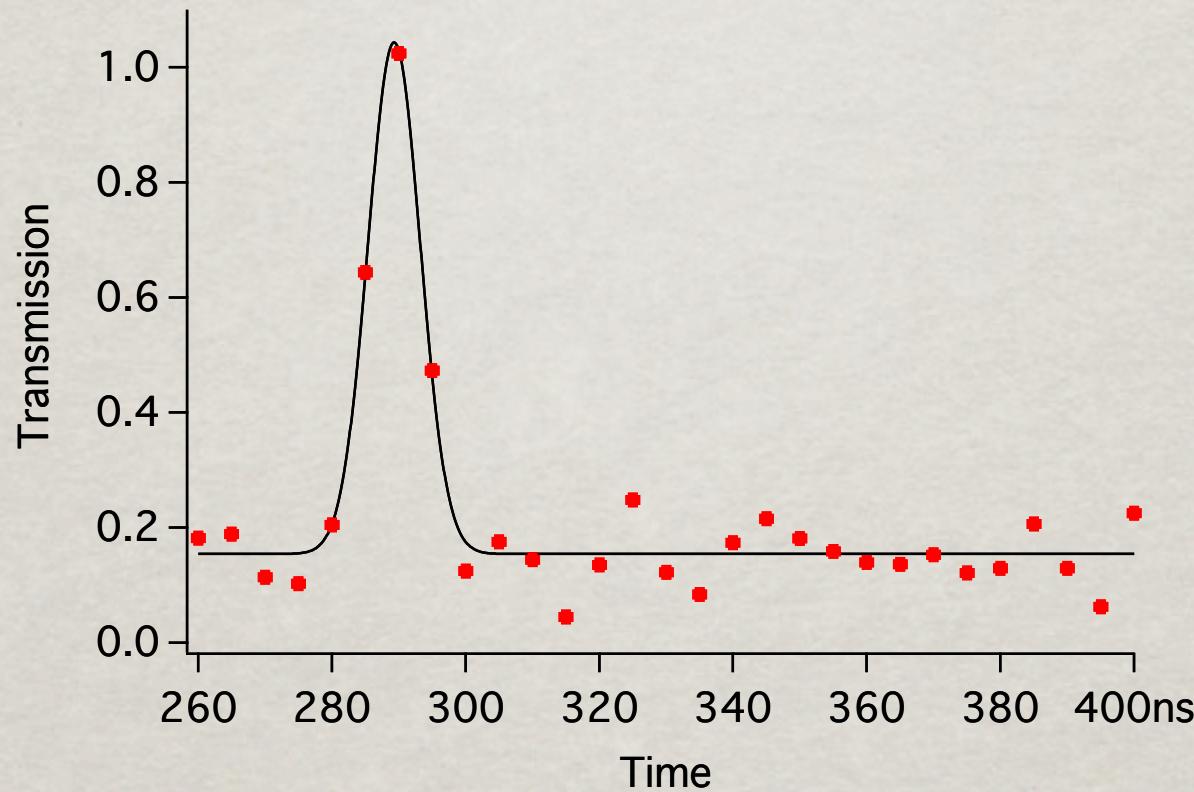


SINGLE-PHOTON ROUTER



- ✿ Operation time down to ~ 10 ns

SINGLE-PHOTON ROUTER

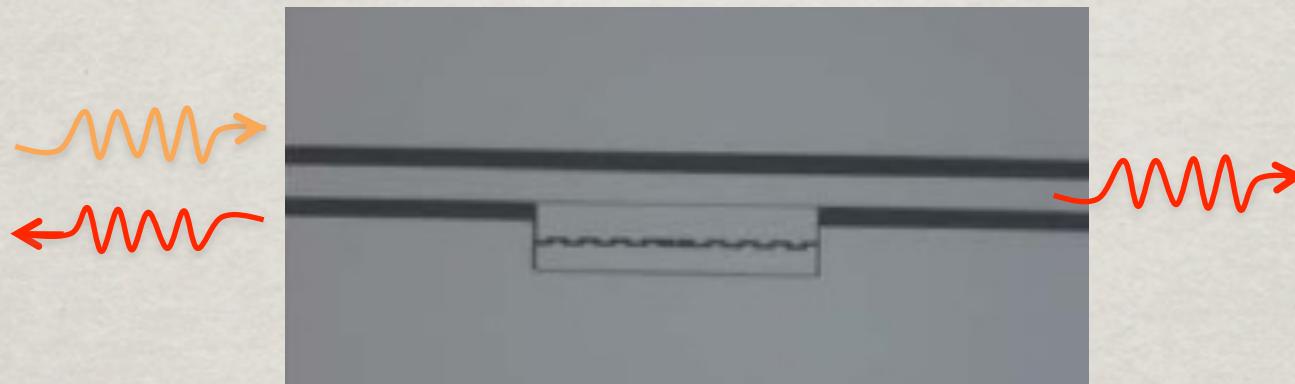


✿ Operation time down to ~ 10 ns

QUANTUM-STATE FILTER

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$$|\Phi_{in}\rangle = a_0 |0\rangle + a_1 |1\rangle + a_2 |2\rangle + \dots$$

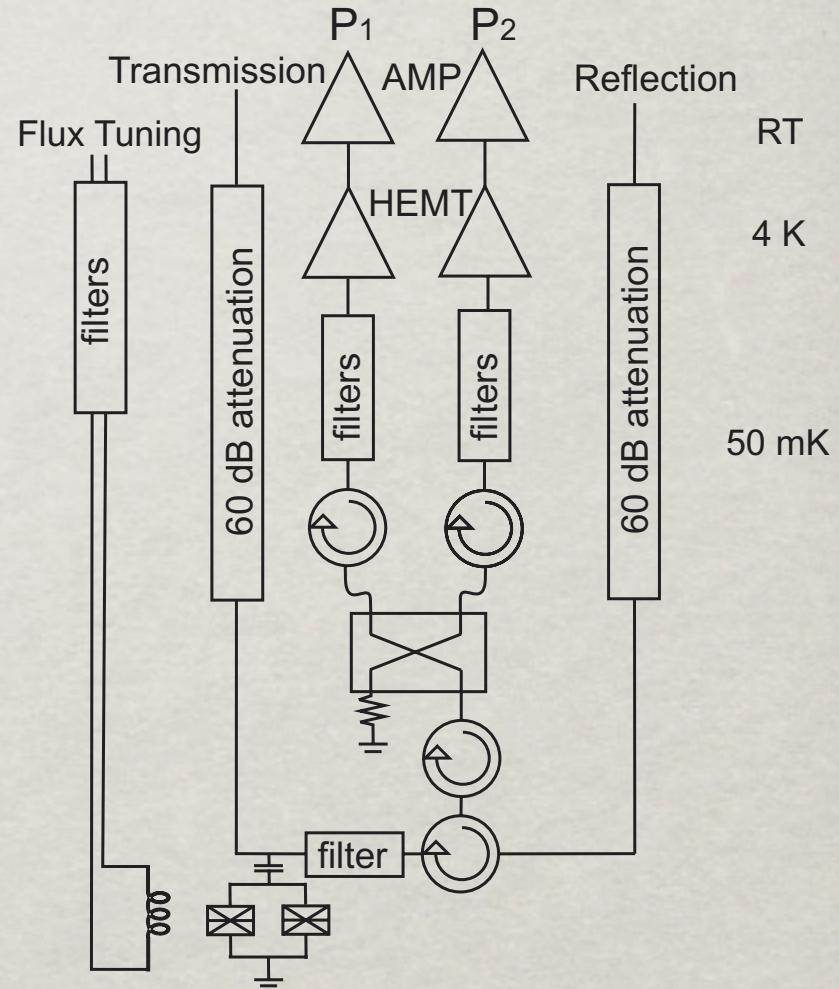


$$|\Phi_R\rangle = r_0 |0\rangle + r_1 |1\rangle \quad |\Phi_T\rangle = t_0 |0\rangle + t_2 |2\rangle + t_3 |3\rangle \dots$$

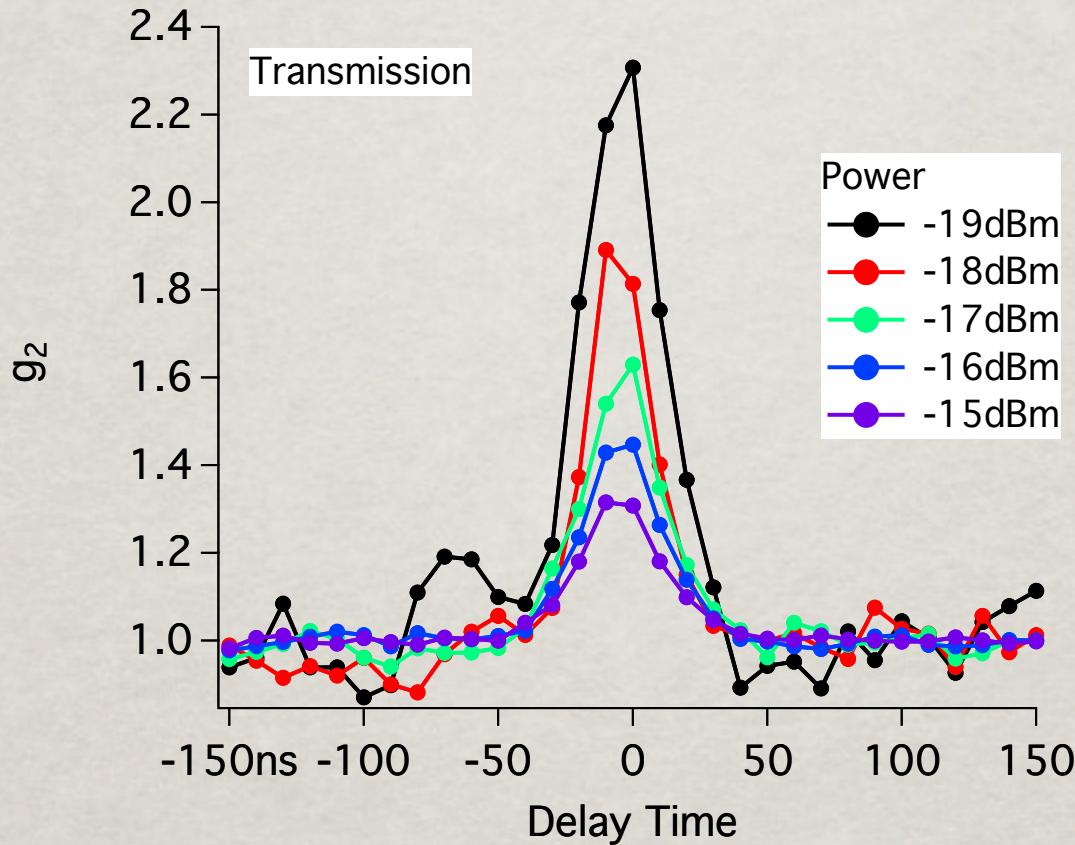
- ✿ Atom selectively reflects 1-photon state
- ✿ Input coherent state converted to nonclassical state

FOURTH MOMENT ($g^{(2)}$)

- HBT measurements of output state
- Commercial “beamsplitter”
- Gabelli et al., PRL (2004)

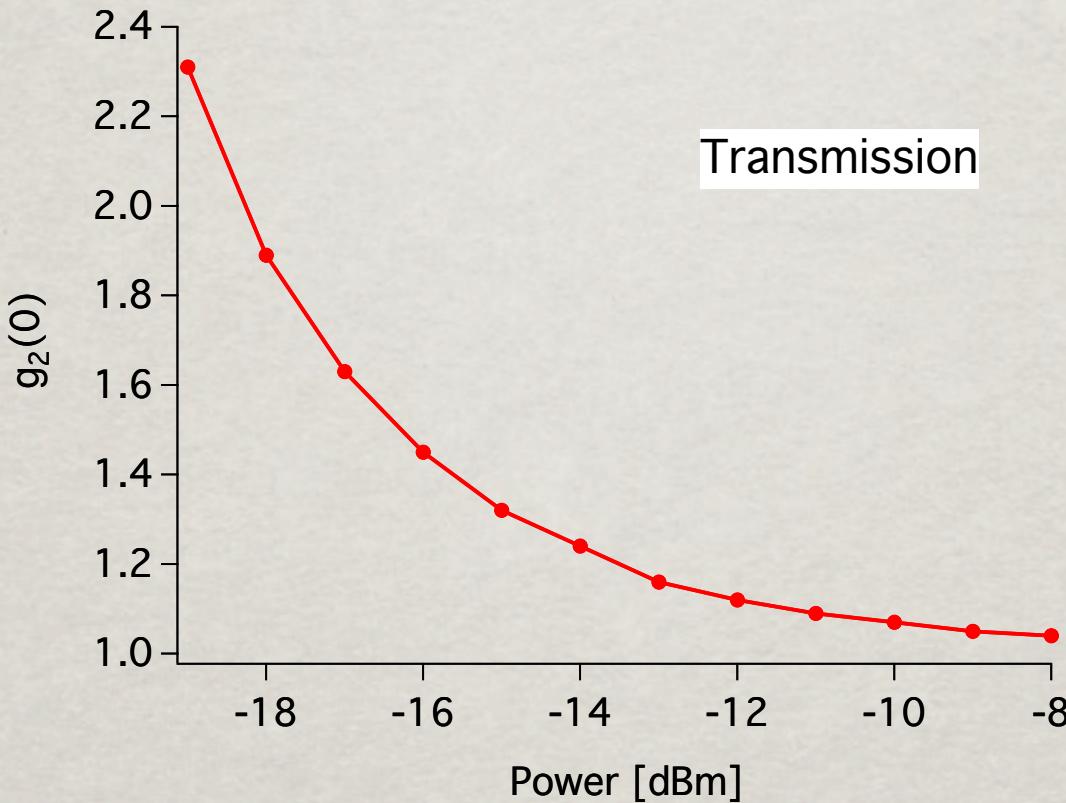


SUPERBUNCHING



- ✿ Observe superbunching of transmitted mode
 - $n = 1$ state “filtered out”

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SUMMARY

- Demonstrated strong extinction caused by a single “atom”
- Demonstrated single-photon router using EIT
- Measured nonclassical statistics produced by a novel quantum-state filter

TOWARDS AN
OPTICAL-MICROWAVE
QUANTUM INTERFACE

HYBRID QUANTUM NETWORK

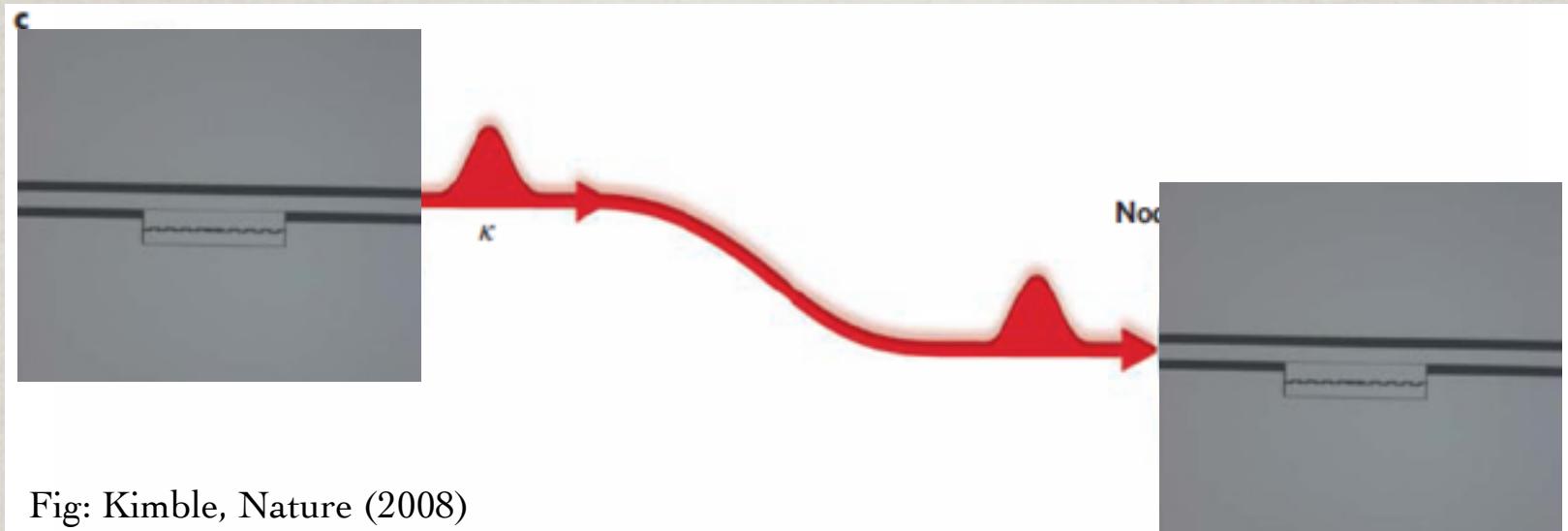
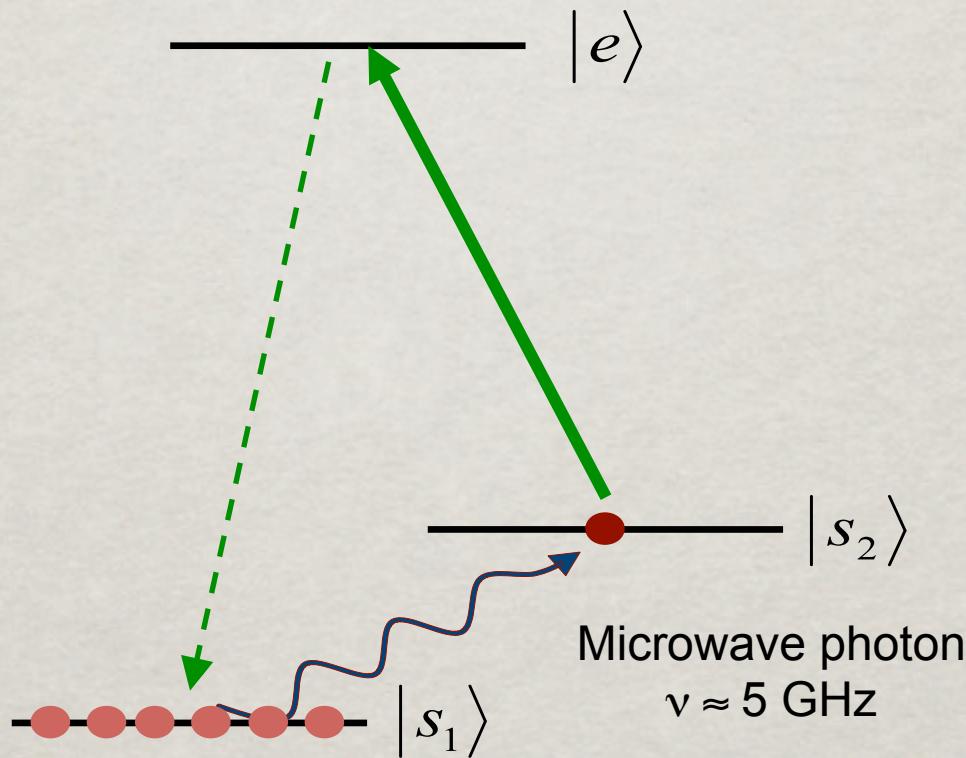


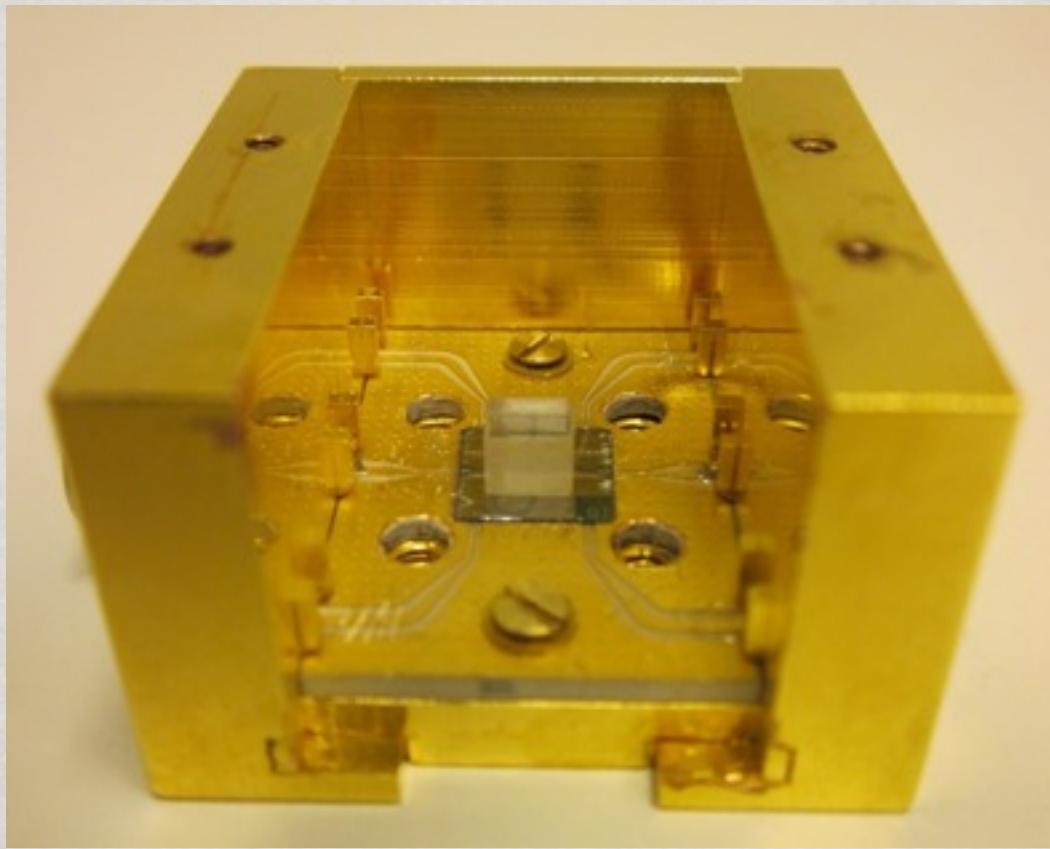
Fig: Kimble, Nature (2008)

- ✿ Telecom photons to distribute quantum information
- ✿ Superconducting circuits to process at nodes
- ✿ Need to coherently convert frequency by factor of 10^4 !

QUANTUM INTERFACE

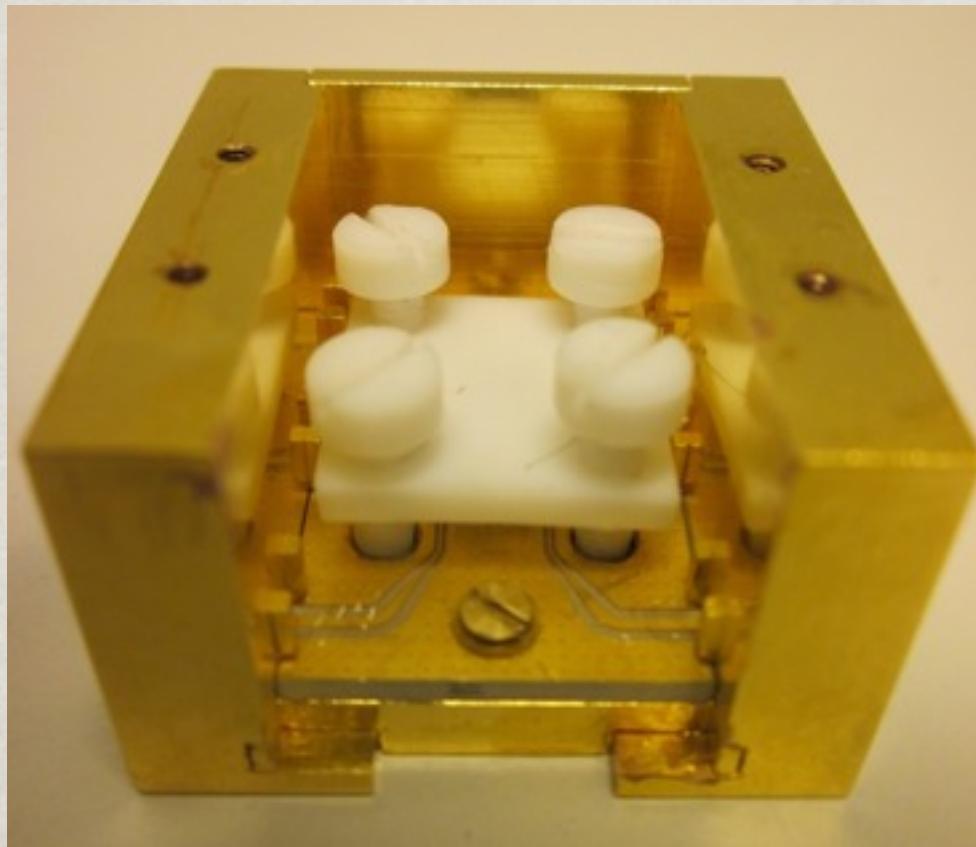


MICROWAVE COUPLING



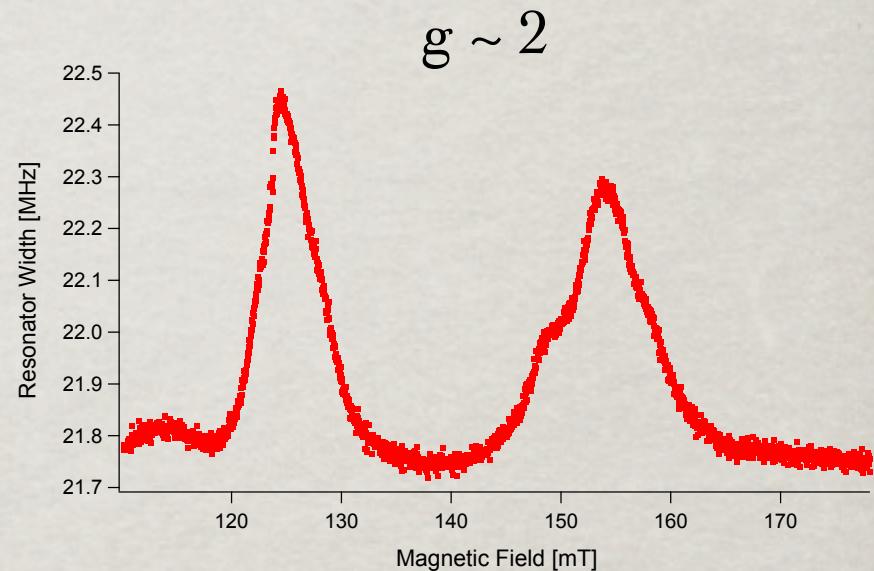
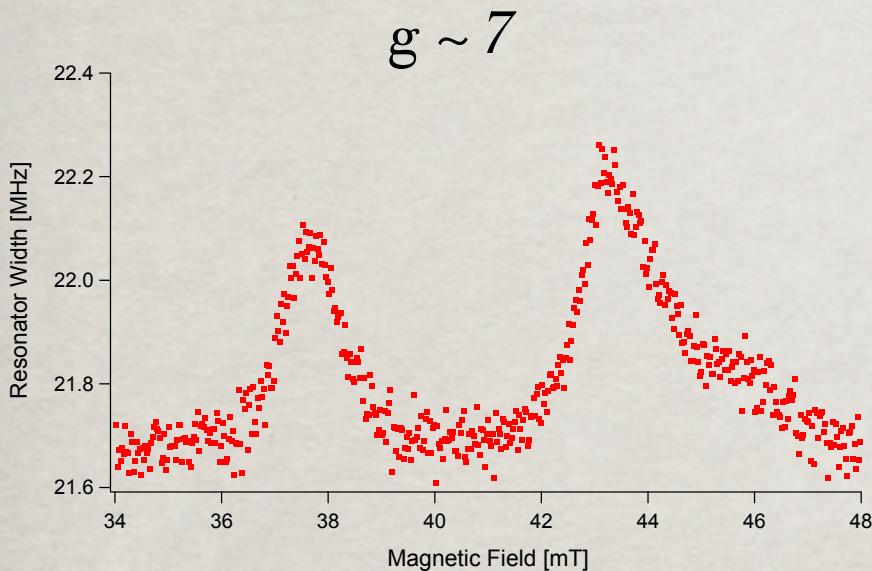
- ✿ Erbium doped crystal coupled to low-Q Nb CPW
- ✿ Pressed contact

MICROWAVE COUPLING



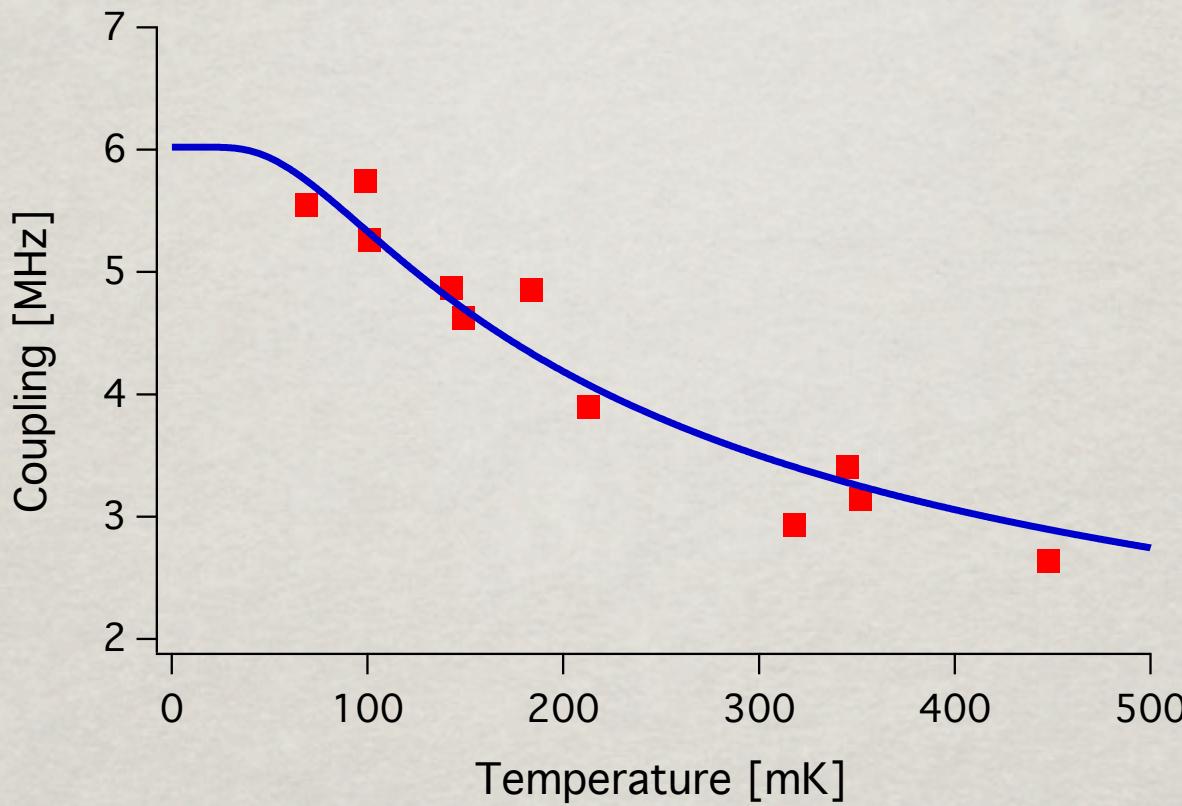
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PLANAR ESR



- ✿ Magnetic field tunes Zeeman splitting of spin levels
- ✿ On resonances, ensemble damps the resonator
- ✿ Measure resonance width vs. magnetic field
- ✿ Collective coupling $g_c \sim 5$ MHz, implying $N \sim 10^{12}$

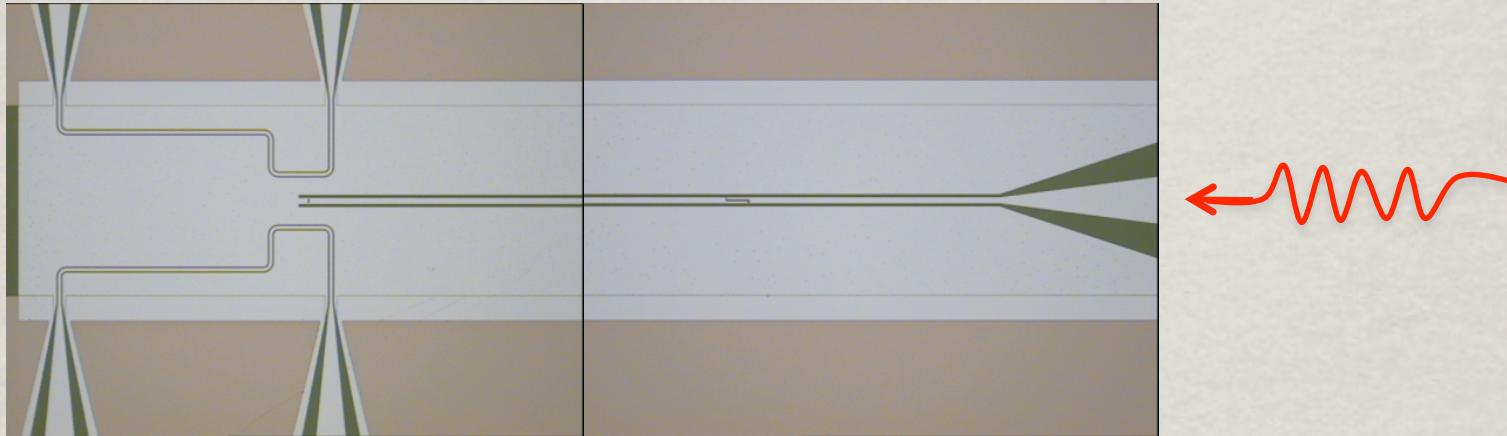
TEMP. DEPENDENCE



- ✿ g_c declines with temp as the spins depolarize

$$g_c = g\sqrt{N(T)} = g_0 \sqrt{\tanh\left(\frac{\hbar\omega}{k_B T}\right)}$$

PROPAGATING PHOTONS



- ❖ Necessary to catch propagating photons originating outside of the cavity
- ❖ “Strong coupling” is not a requirement

FUTURE WORK

- ✿ Tunable resonator to measure zero field lines
- ✿ Compare with Nd
- ✿ Impedance matching and photonic memory...